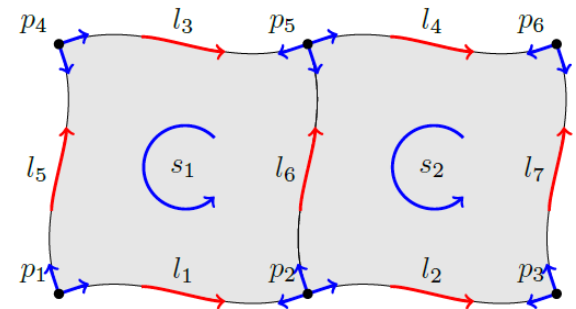
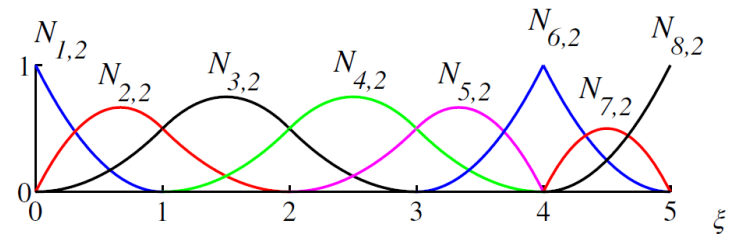


Midterm review:

Mimetic Isogeometric FEM

M.Sc. Thesis project
by Stevie-Ray Janssen

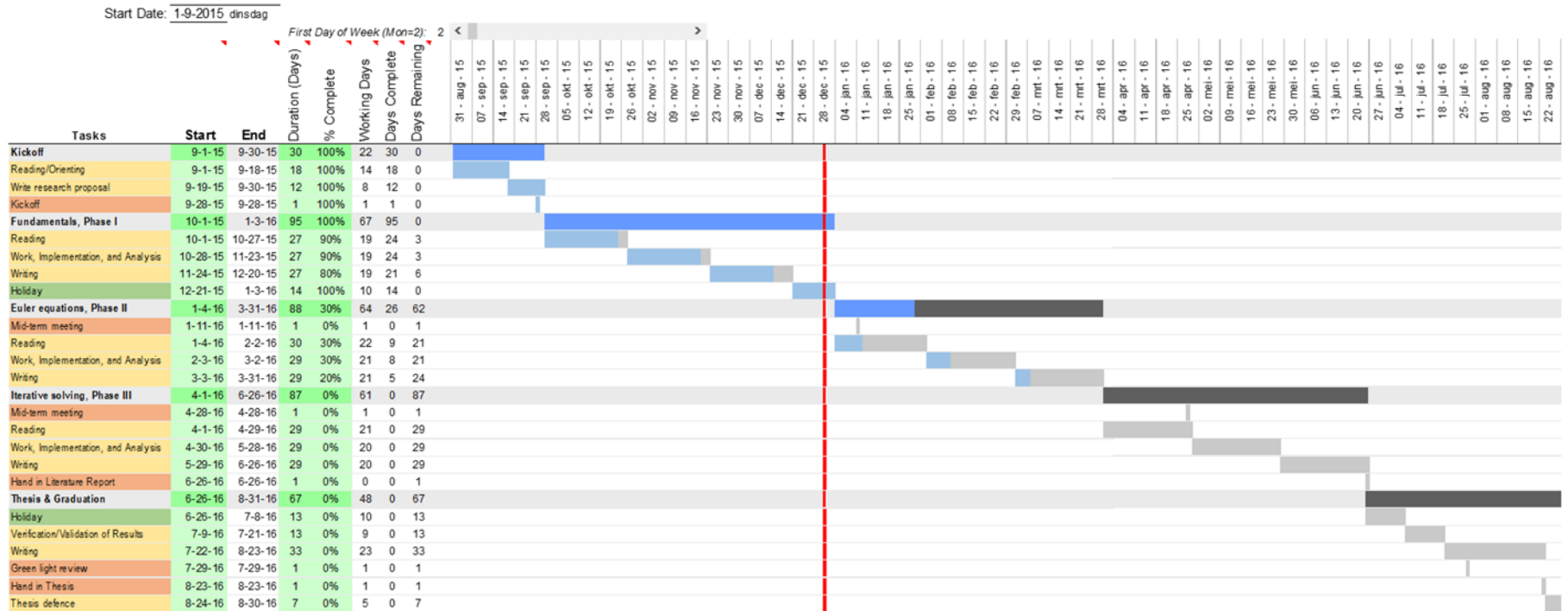


Objective:

Combine ideas from isogeometric analysis and mimetic methods to develop a structure-preserving discretization for the Euler equations for incompressible fluids.

Project outline

- Planning:



Project outline (cnt'd)

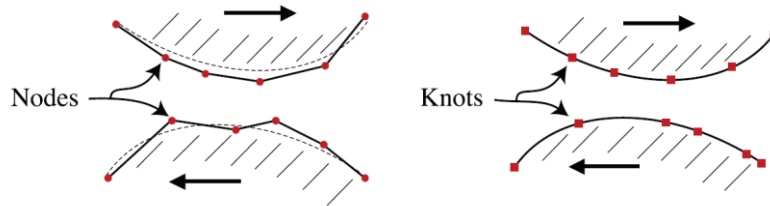
- Phase I questions:
 - How can we use IGA to solve PDE's?
 - What structures are facilitated in elliptic PDE's?
 - How can we preserve these structures?
 - Can we construct a MIMIGA method to discretize an elliptic PDE problem?

This presentation - literature review

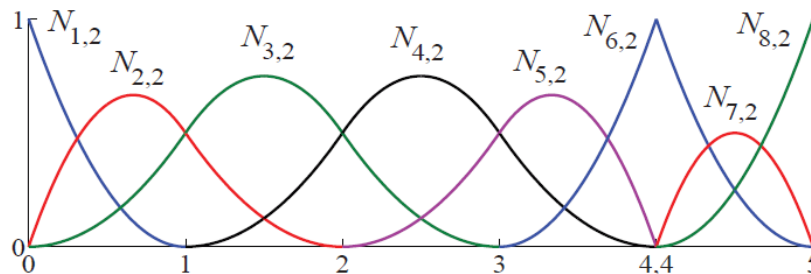
- Introduction
 - Isogeometric Analysis & Mimetic Methods
- Approach for elliptic PDE's
 - Exterior calculus
 - DeRham complex
 - Application: Scalar Poisson equation in 2D
- Conclusion
- Future work

Introduction – Isogeometric Analysis

- Introduced by the Hughes group in 2005 to bridge the gap between CAD and FEM
- Isogeometric paradigm



- B-splines make an excellent basis for FEM



Introduction – Mimetic Methods

- PDE's facilitate physical structures and symmetries.
- Tools from exterior calculus and algebraic topology are used to capture these structures.
- Growing awareness: Discrete exterior calculus, discrete hodge theory, exterior finite element method, compatible methods, mimetic finite difference, etc

Why exterior calculus?

- Structures become apparent.
- Distinction between topological and metric dependencies.
- Generalized for n dimensions.

Differential Forms; $\alpha^{(k)}$

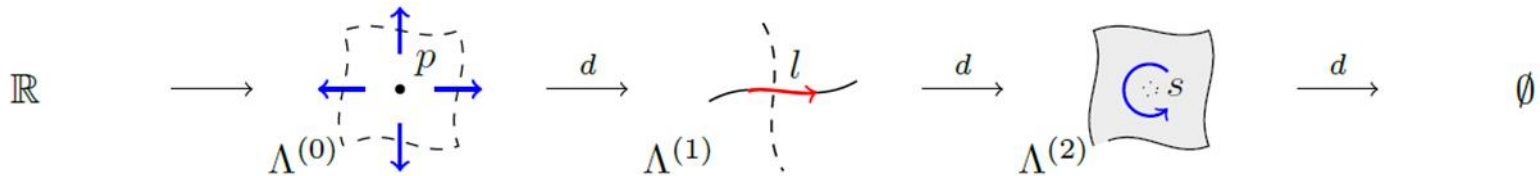
- Differential forms are elements from the dual vector space,
- Associated with geometric structure,
 - 0-form: $f^{(0)} = f(x, y)$
 - 1-form: $\alpha^{(1)} = \alpha_1(x, y)dx + \alpha_2(x, y)dy$
- “Measurement of physical variables,”
 - $M = \iint \rho^{(2)} = \iint \rho(x, y)dx \wedge dy$
- Space of k-forms: $\Lambda^{(k)}$

Exterior derivative; d

- Exterior derivative d generalizes ∇f , $\nabla \times \underline{\omega}$, $\nabla \cdot \underline{v}$

$$d\alpha^{(1)} = \left(\frac{\partial \alpha_2}{\partial x} - \frac{\partial \alpha_1}{\partial y} \right) dx \wedge dy$$

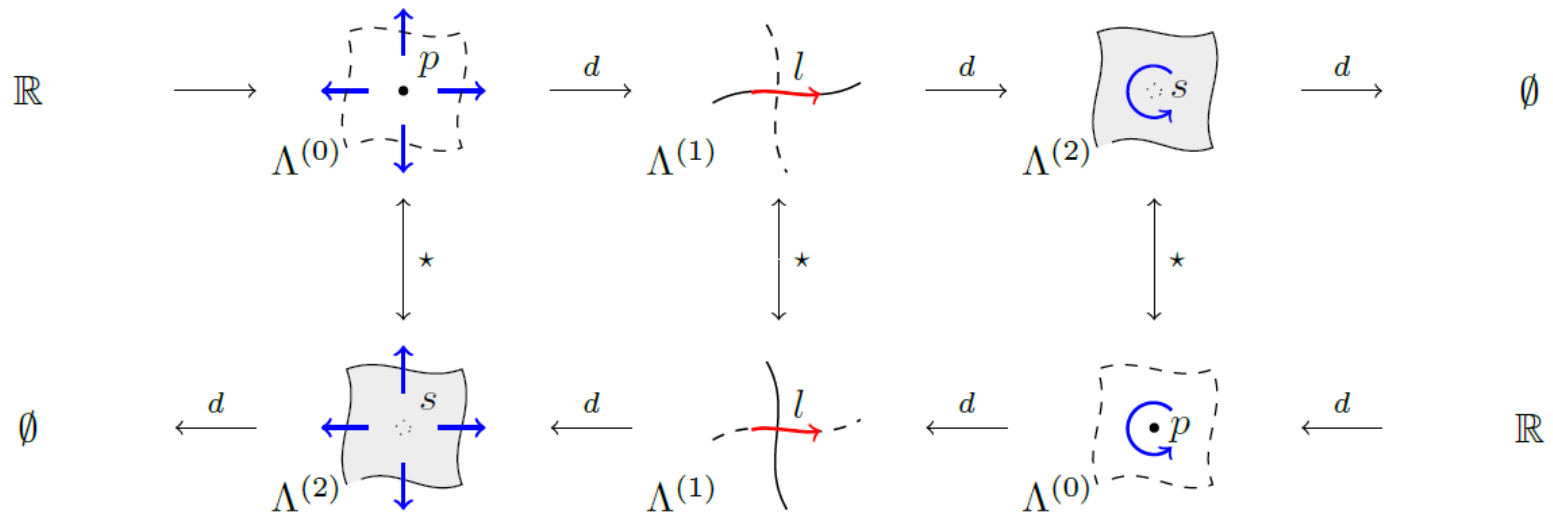
- $d: \Lambda^{(k)} \rightarrow \Lambda^{(k+1)}$



- Exact sequence, the DeRham complex
- Nilpotent, $dd\alpha^{(k)} = 0$
- Independent of metric

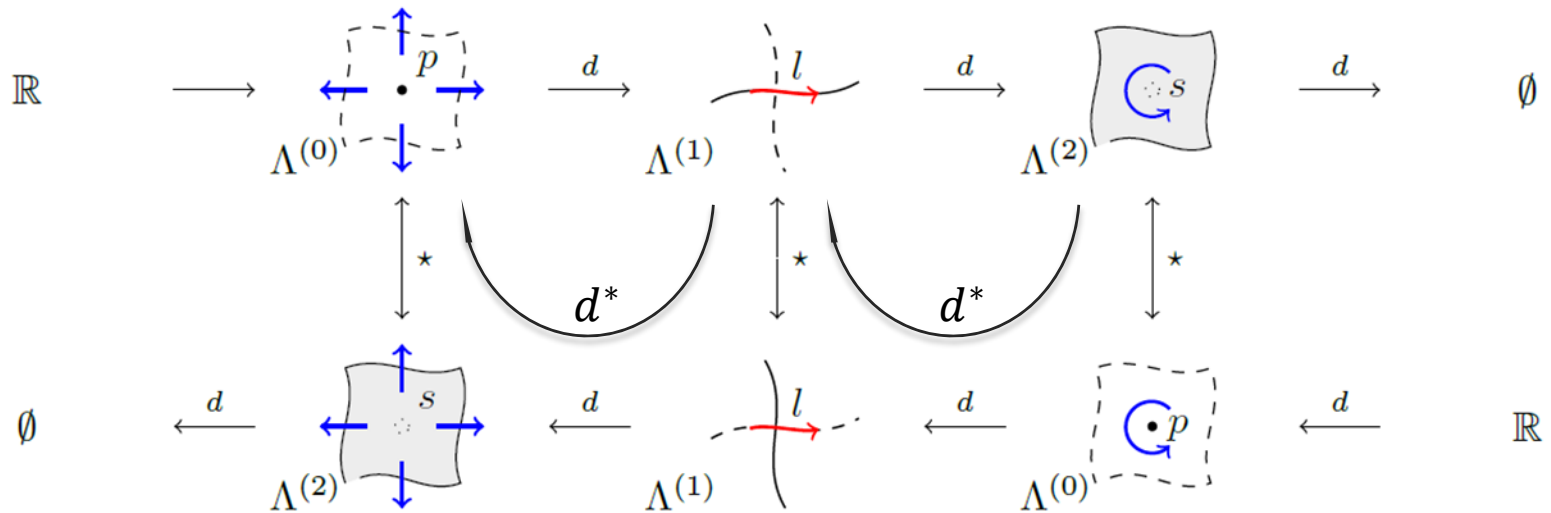
Hodge- \star operator;

- Maps forms to dual geometry,
- Metric dependent,
- Double DeRham complex,

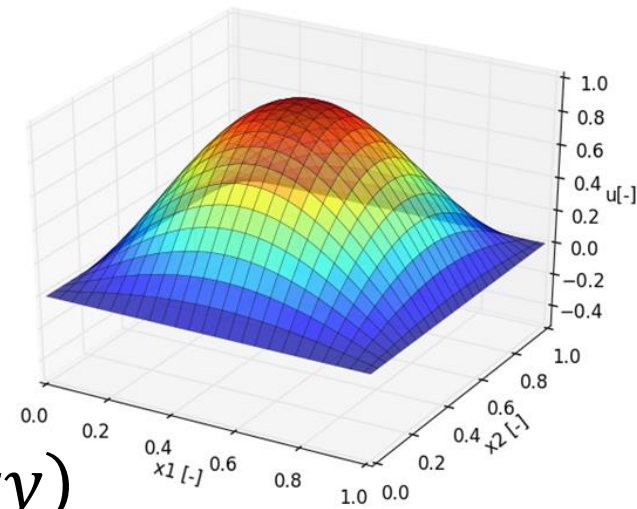


Codifferential; d^*

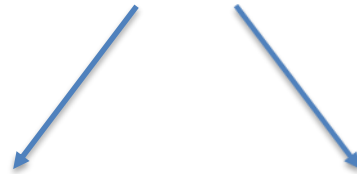
- $d^* := \star d \star$
- Adjoint of d : $(\cdot, d^* \cdot) = (d \cdot, \cdot) - \int \text{bc's}$
- Laplace operator: $\Delta = dd^* + d^*d$



Scalar Poisson equation



- E.g. Potential flow, electrostatics,
- Given $f(x, y) = 2\pi^2 \sin(\pi x) \sin(\pi y)$
find $\varphi(x, y)$ such that $\Delta\varphi = f$ on $\Omega = [0, 1]^2$
with $\varphi = 0$ on $\partial\Omega$



0-form,

Find $\varphi^{(0)}$ s.t. $d^*d\varphi^{(0)} = f^{(0)}$

2-form,

Find $\sigma^{(2)}$ s.t. $dd^*\sigma^{(2)} = f^{(2)}$

Same solution, different discretization

0-form Poisson; $d^* d\varphi^{(0)} = f^{(0)}$

- Weak formulation,

$$(w^{(0)}, d^* d\varphi^{(0)})_{\Omega} = (w^{(0)}, f^{(0)})_{\Omega}$$

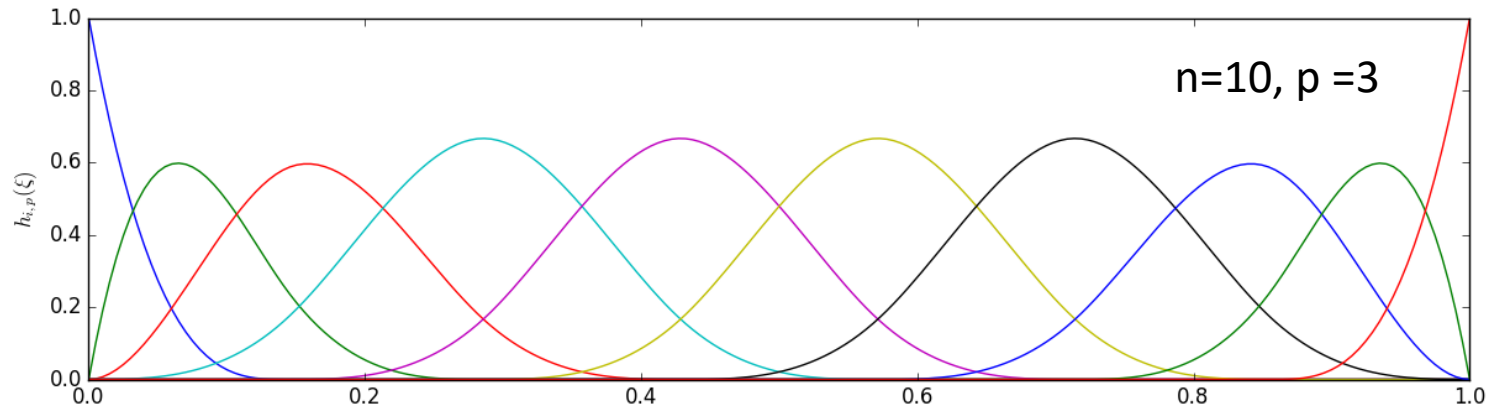
\Leftrightarrow

$$(dw^{(0)}, d\varphi^{(0)})_{\Omega} = (w^{(0)}, f^{(0)})_{\Omega} - \int_{\partial\Omega} w^{(0)} \Lambda^* d\varphi^{(0)}$$

- Well-posedness through Lax-Milgram,

0-form Poisson; FEM

- Conforming FEM, take $\Lambda_h^{(k)} \subset \Lambda^{(k)}$
- Use B-spline spans $\Lambda_h^{(0)} = S^{p,p}$



0-form Poisson; edge functions

- Applying the exterior derivative (1D-example)

- Nodal basis: $\varphi_h^{(0)} = \sum_{i=0}^n \varphi_i h_i^p(x) = (\underline{\varphi})^T \underline{R}^0$

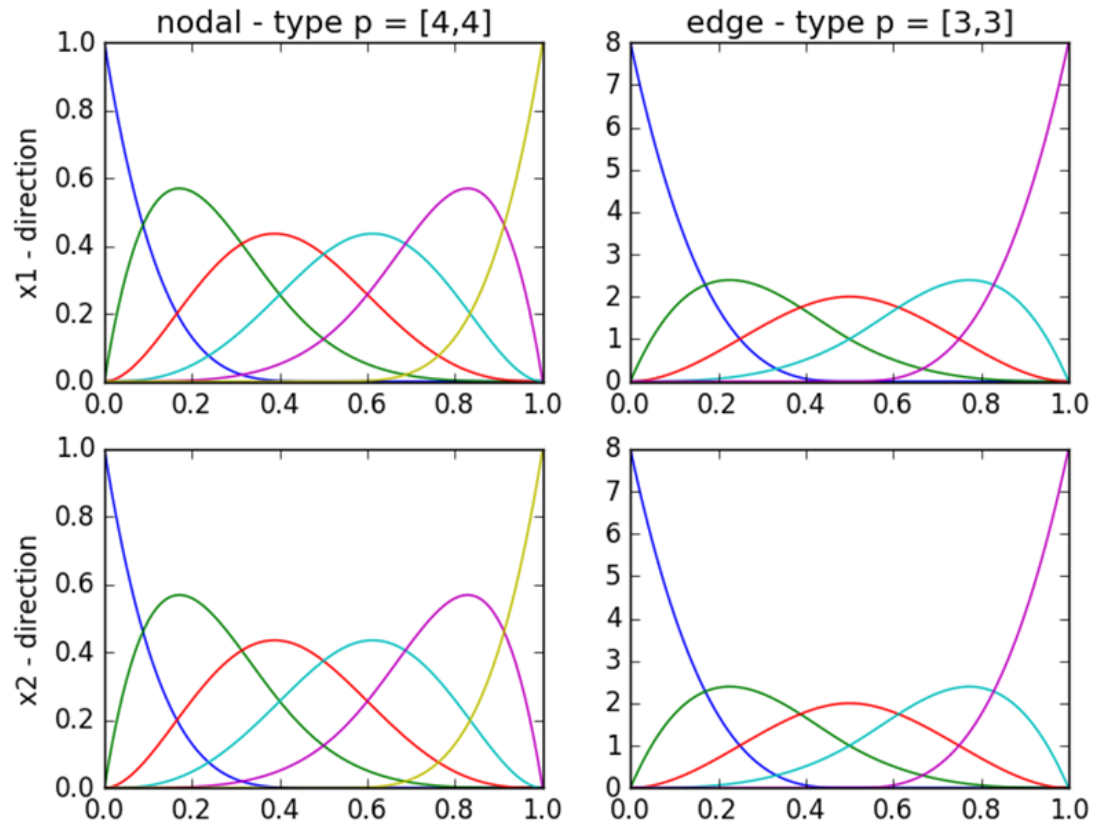
- Then, $d\varphi_h^{(0)} = \sum_{i=1}^n (\varphi_i - \varphi_{i-1}) e_i^{p-1}(x) = (\mathbb{E}^{(10)} \underline{\varphi})^T \underline{R}^1$

Differences of coefficients
are captured in matrix
using $\{-1,0,1\}$

New edge type basis function
emerges with a polynomial
degree less


0-form Poisson; edge functions (cnt'd)

- Extension to 2D using tensor products of nodal and edge type basis
- Nodal/edge
 - 0-form
 - 1-form
 - 2-form



0-form Poisson, Matrices

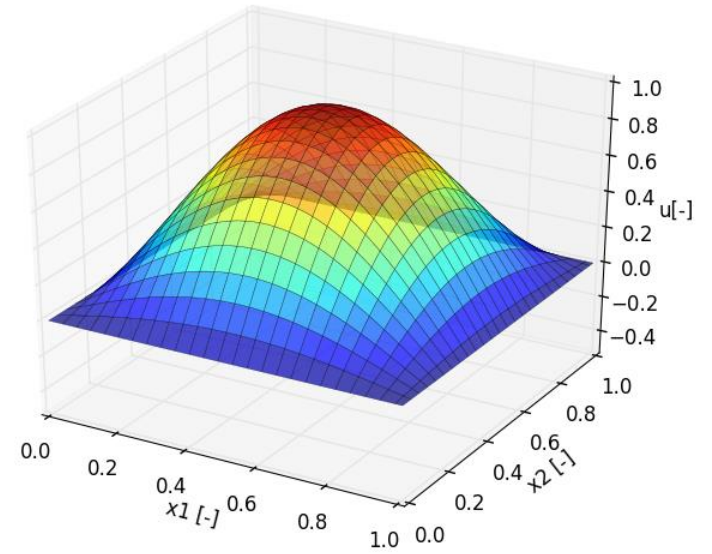
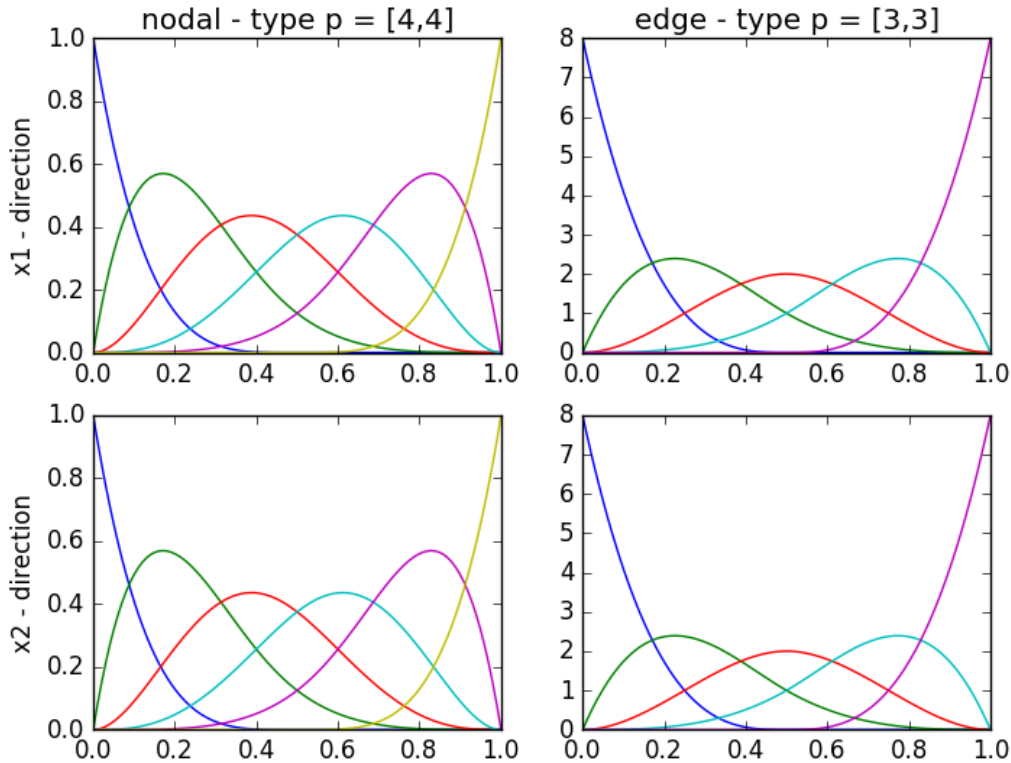
- $\left(dw_h^{(0)}, d\varphi_h^{(0)} \right)_\Omega = \underline{w}^T (\mathbb{E}^{10})^T \left(\int_\Omega (\underline{R}^{(1)})^T \underline{R}^{(1)} \right) (\mathbb{E}^{10}) \underline{\varphi}$

Mass matrix $\mathbb{M}^{(11)}$

- Result: $(\mathbb{E}^{(10)})^T \mathbb{M}^{(11)} \mathbb{E}^{(10)} \underline{\varphi} = \underline{f}$

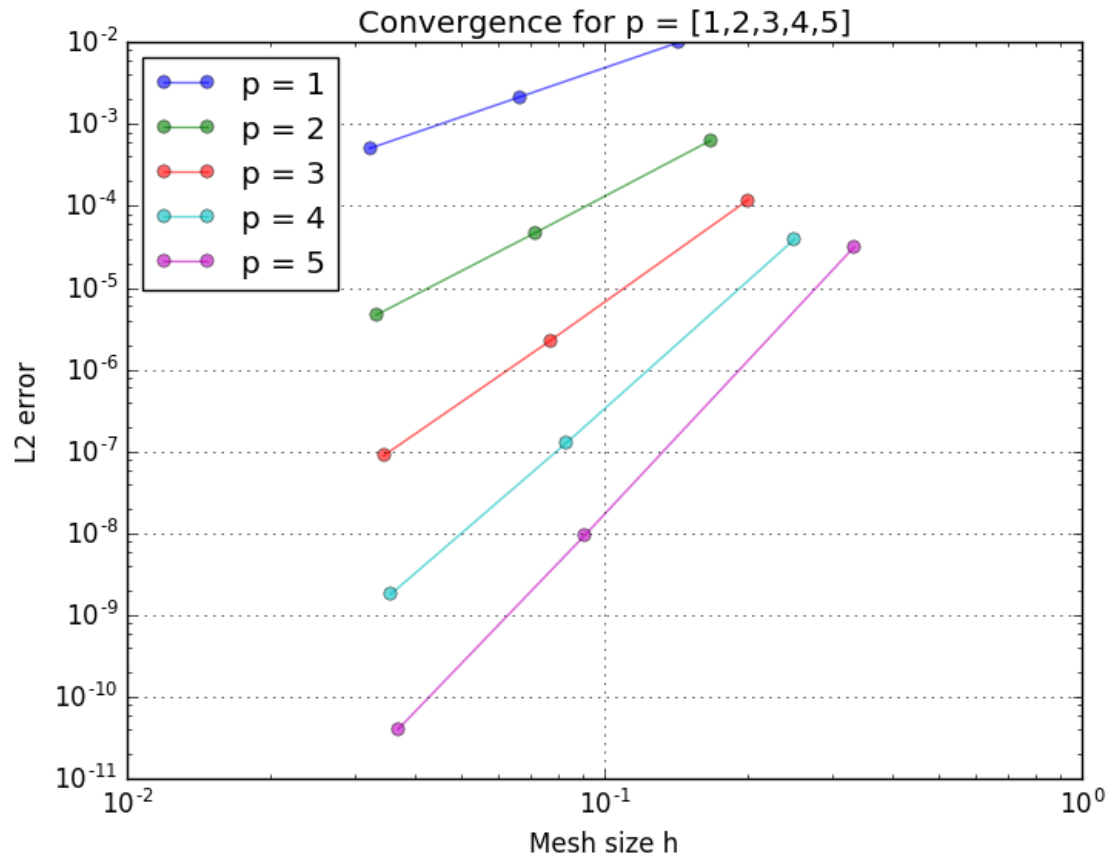
0-form Poisson, Matrices (cnt'd)

- Exact discretization of $v^{(1)} = d\varphi^{(0)}$ through incidence matrices, $\underline{v} = \mathbb{E}^{(10)} \underline{\varphi}$
- Incidence matrices are nilpotent $\mathbb{E}^{(21)} \mathbb{E}^{(10)} = \emptyset$, and satisfy the DeRham sequence
- Hodge- \star operator (metric) is discretized through mass matrix $\mathbb{M}^{(11)}$

0-form Poisson, Results



0-form Poisson, Results (cnt'd)



P	Slope
1	1.9998
2	3.0444
3	4.0945
4	5.1311
5	6.1736

2-form Poisson; $dd^* \sigma^{(2)} = f^{(2)}$

- Weak formulation;

$$(w^{(2)}, dd^* \sigma^{(2)})_{\Omega} = (w^{(2)}, f^{(2)})_{\Omega}$$

- Integration by parts? No, take mixed formulation:

$$\begin{cases} d^* \sigma^{(2)} = \psi^{(1)} \\ d\psi^{(1)} = f^{(2)} \end{cases}$$

- Weak form:

$$\begin{cases} (dq^{(1)}, \sigma^{(2)})_{\Omega} = (q^{(1)}, \psi^{(1)})_{\Omega} - \int_{\partial\Omega} q^{(1)} \wedge \star \sigma^{(2)} \\ (w^{(2)}, d\psi^{(1)})_{\Omega} = (w^{(2)}, f^{(2)})_{\Omega} \end{cases}$$

- Well posedness through Inf-Sup conditions

2-form Poisson; FEM

- Can we take,
 - $\Lambda_h^{(1)} = S^{p,p}$?
 - $\Lambda_h^{(2)} = S^{p,p}$?
- No, well-posedness depends on the DeRham sequence. We take
 - $\Lambda_h^{(1)} = S^{p-1,p} \times S^{p,p-1}$
 - $\Lambda_h^{(2)} = S^{p-1,p-1}$
- Which satisfy exact sequence

$$S^{p,p} \xrightarrow[\mathfrak{d}]{\mathbb{E}^{(10)}} S^{p-1,p} \times S^{p,p-1} \xrightarrow[\mathfrak{d}]{\mathbb{E}^{(21)}} S^{p-1,p-1}$$

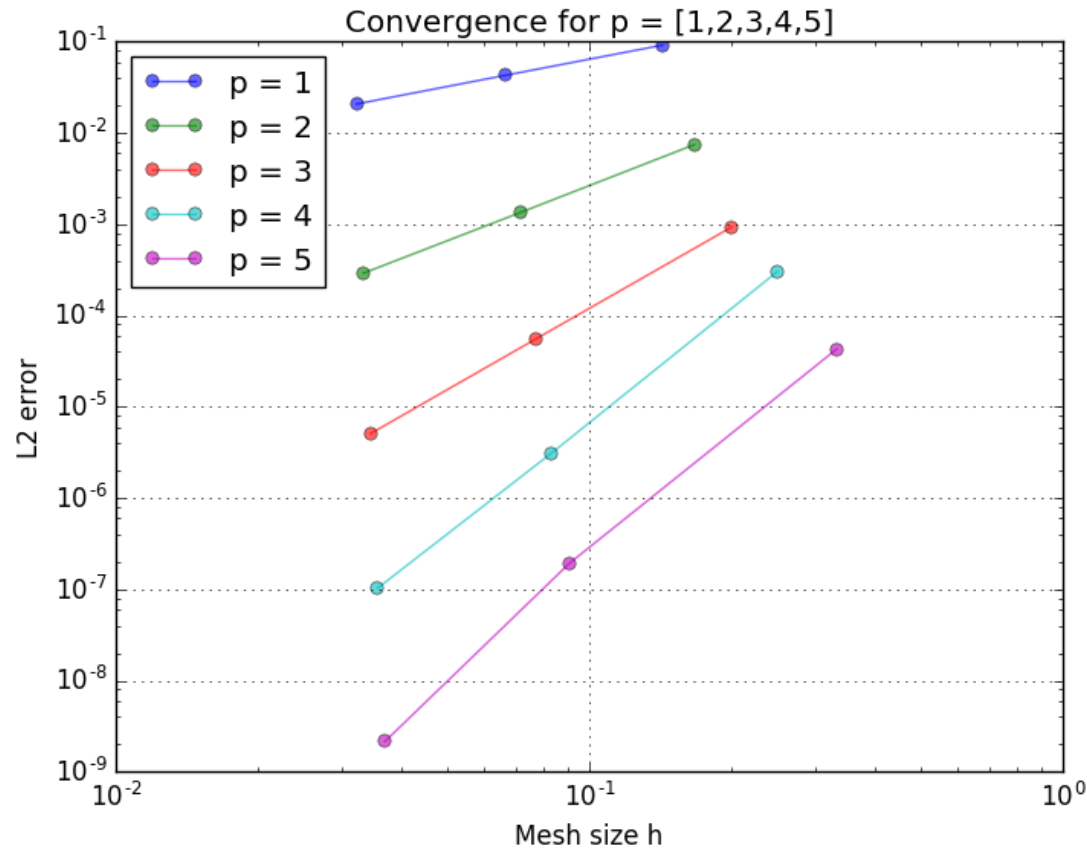
2-form Poisson; Matrices

$$\begin{cases} -(q^{(1)}, \psi^{(1)})_{\Omega} + (dq^{(1)}, \sigma^{(2)})_{\Omega} = 0 \\ (w^{(2)}, d\psi^{(1)})_{\Omega} = (w^{(2)}, f^{(2)})_{\Omega} \end{cases}$$

$$\begin{bmatrix} -\mathbb{M}^{(11)} & (\mathbb{M}^{(22)} \mathbb{E}^{(21)})^T \\ \mathbb{M}^{(22)} \mathbb{E}^{(21)} & \emptyset \end{bmatrix} \begin{bmatrix} \underline{\psi} \\ \underline{\sigma} \end{bmatrix} = \begin{bmatrix} \underline{0} \\ \underline{f} \end{bmatrix}$$

- Or $(\mathbb{M}^{(22)} \mathbb{E}^{(21)})^T (\mathbb{M}^{(11)})^{-1} (\mathbb{M}^{(22)} \mathbb{E}^{(21)}) \underline{\psi} = \underline{f}$

2-form Poisson; Results



P	Slope
1	0.9977
2	2.0150
3	2.9672
4	4.1104
5	4.4811

Conclusion

- Elliptic problems can be discretized using mass matrices and incidence matrices.
- Solution spaces are chosen such that they satisfy the DeRham complex.

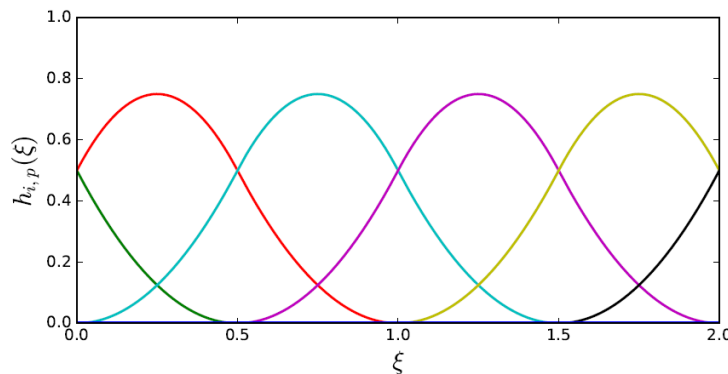
Conclusion (cnt'd)

- Comparison 0-form & 2-form Poisson:

0-form	2-form
$(\mathbb{E}^{(10)})^T \mathbb{M}^{(11)} \mathbb{E}^{(10)}$	$\begin{bmatrix} -\mathbb{M}^{(11)} & (\mathbb{M}^{(22)} \mathbb{E}^{(21)})^T \\ \mathbb{M}^{(22)} \mathbb{E}^{(21)} & \emptyset \end{bmatrix}$
Obtain solution $\varphi^{(0)}$	Obtain solutions $\sigma^{(2)}, \psi^{(1)}$
<ul style="list-style-type: none"> - Dirichlet is essential - Neumann is natural 	<ul style="list-style-type: none"> - Dirichlet is natural - Neumann is essential
Gradient exact $\mathbb{E}^{(10)} \underline{\varphi} = \underline{v}$	Divergence exact $\mathbb{E}^{(21)} \underline{\psi} = \underline{0}$ i.e. $\nabla \cdot v = 0$

Future Work

- Towards the incompressible Euler equations:
 - Extend to hyperbolic problems,
 - Linear advection equation.
 - Construction of periodic domain.



- Staggering velocity and vorticity in time?

Questions