

Modelling the interaction between the electromagnetic field and fluids

Literature review

Michiel de Reus

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Introduction

- Electromagnetic field: Maxwell equations
- Fluid dynamics: Navier-Stokes equations
- Coupling through dependent terms:
 - Induced currents and constitutive relations depend on velocity field;
 - External force and source term depend on EM fields.

Maxwell equations

Maxwell equations in matter:

$$(1) \quad -\nabla \times \mathbf{H} + \frac{\partial \mathbf{D}}{\partial t} = -\mathbf{J}^f,$$

$$(2) \quad \nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0,$$

$$(3) \quad \nabla \cdot \mathbf{D} = \rho_f,$$

$$(4) \quad \nabla \cdot \mathbf{B} = 0.$$

Maxwell equations

Fields:

- $\mathbf{E}(\mathbf{x}, t)$, electric field;
- $\mathbf{H}(\mathbf{x}, t)$, magnetic field;
- $\mathbf{D}(\mathbf{x}, t)$, electric flux field;
- $\mathbf{B}(\mathbf{x}, t)$, magnetic flux field.

Other quantities:

- $\mathbf{J}^f(\mathbf{x}, t)$, free current density,
- $\rho_f(\mathbf{x}, t)$, free charge density.

Constitutive relations

We need relations between the different fields, constitutive relations.

For non-dispersive linear conductive media:

- $\mathbf{D} = \varepsilon \mathbf{E}$,
- $\mathbf{H} = \mu \mathbf{B}$.
- $\rho_f = 0$,
- $\mathbf{J}^f = \sigma \mathbf{E}$.

Constitutive relations

Up until now no flow. Assume velocity field \mathbf{v} , then using Lorentz invariance:

$$(5) \quad \mathbf{D} = \epsilon \mathbf{E} + \epsilon \mathbf{v} \times \mathbf{B} - \frac{1}{c^2} \mathbf{v} \times \mathbf{H},$$

$$(6) \quad \mathbf{B} = \mu \mathbf{H} - \mu \mathbf{v} \times \mathbf{D} + \frac{1}{c^2} \mathbf{v} \times \mathbf{E}.$$

Constitutive relations

Linearized and solved for \mathbf{E} and \mathbf{H} :

$$(7) \quad \mathbf{D} \approx \varepsilon \mathbf{E} + (\mu\varepsilon - \mu_0\varepsilon_0) \mathbf{v} \times \mathbf{H},$$

$$(8) \quad \mathbf{B} \approx \mu \mathbf{H} - (\mu\varepsilon - \mu_0\varepsilon_0) \mathbf{v} \times \mathbf{E}.$$

Note:

$$\mu\varepsilon - \mu_0\varepsilon_0 = \frac{1}{c_m^2} - \frac{1}{c^2}.$$

Constitutive relations

For electric charge and current:

$$\rho_f = \gamma \frac{1}{c^2} \sigma \mathbf{E} \cdot \mathbf{v},$$
$$\mathbf{J}^f = \gamma \sigma \mathbf{E} + \gamma \sigma \mathbf{v} \times \mathbf{B}.$$

If $v \ll c$ and $\sigma \ll 1$:

$$\rho_f \approx 0,$$
$$\mathbf{J}^f \approx \sigma \mathbf{E} + \sigma \mathbf{v} \times \mathbf{B}.$$

Maxwell equations

In terms of E and H , with external sources:

$$(9) \quad -\nabla \times \mathbf{H} + \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} = -\mathbf{J}^{\text{ind}} - \mathbf{J}^{\text{ext}},$$

$$(10) \quad \nabla \times \mathbf{E} + \mu_0 \frac{\partial \mathbf{H}}{\partial t} = -\mathbf{K}^{\text{ind}} - \mathbf{K}^{\text{ext}}.$$

- Induced currents are functions of the fields,
- External currents are independent of the field (e.g. the lab laser).

EM energy in vacuum

$$(11) \quad \partial_t u_{em} = -\nabla \cdot \mathbf{S} - \mathbf{E} \cdot \mathbf{J} - \mathbf{H} \cdot \mathbf{K},$$

EM energy density

$$(12) \quad u_{em} = \frac{1}{2} (\epsilon_0 \|\mathbf{E}\|^2 + \mu_0 \|\mathbf{H}\|^2),$$

Poynting vector

$$(13) \quad \mathbf{S} = \mathbf{E} \times \mathbf{H}.$$

Lorentz force

Regular form for point charges:

$$\mathbf{F} = q [\mathbf{E} + \mu_0 \mathbf{v} \times \mathbf{H}] .$$

Continuous form

$$\mathbf{f} = \rho_e \mathbf{E} + \mu_0 \mathbf{J} \times \mathbf{H},$$

with electric charge density

$$\rho_e = \varepsilon_0 \nabla \cdot \mathbf{E}.$$

Lorentz force

Generalization with magnetic currents:

$$\mathbf{f} = \rho_e \mathbf{E} + \rho_m \mathbf{H} + \mu_0 \mathbf{J} \times \mathbf{H} - \varepsilon_0 \mathbf{K} \times \mathbf{E}.$$

with magnetic charge density

$$\rho_m = \mu_0 \nabla \cdot \mathbf{H}.$$

From Maxwell equations we can derive:

$$\mathbf{f} = \nabla \cdot \mathbf{T} - \frac{\partial \mathbf{S}}{\partial t},$$

Stress tensor

Stress tensor in components:

$$(14) \quad T_{ij} = \mu_0 H_i H_j + \varepsilon_0 E_i E_j - \frac{1}{2} \delta_{ij} [\varepsilon_0 E_i E_i + \mu_0 H_i H_i].$$

In general we need 12 quantities.

Currents without velocity

For linear conducting media:

$$(15) \quad \mathbf{J}^{\text{ind}} = (\varepsilon - \varepsilon_0) \frac{\partial \mathbf{E}}{\partial t} + \sigma \mathbf{E},$$

$$(16) \quad \mathbf{K}^{\text{ind}} = (\mu - \mu_0) \frac{\partial \mathbf{H}}{\partial t}.$$

- External magnetic currents zero.
- External electric current depend on situation.

EM energy in (linear) matter

$$(17) \quad \partial_t u_{em} = -\nabla \cdot \mathbf{S} - \sigma \|\mathbf{E}\|^2 - \mathbf{E} \cdot \mathbf{J}^{ext},$$

EM energy density

$$(18) \quad u_{em} = \frac{1}{2} (\epsilon \|\mathbf{E}\|^2 + \mu_0 \|\mathbf{H}\|^2),$$

Poynting vector

$$(19) \quad \mathbf{S} = \mathbf{E} \times \mathbf{H}.$$

Currents with velocity

For linear conducting media:

$$\begin{aligned}\mathbf{J}^{\text{ind}} &= (\varepsilon - \varepsilon_0) \frac{\partial \mathbf{E}}{\partial t} + (\mu\varepsilon - \mu_0\varepsilon_0) \frac{\partial}{\partial t} (\mathbf{v} \times \mathbf{H}) \\ &\quad + \sigma (\mathbf{E} + \mu\mathbf{v} \times \mathbf{H}), \\ \mathbf{K}^{\text{ind}} &= (\mu - \mu_0) \frac{\partial \mathbf{H}}{\partial t} - (\mu\varepsilon - \mu_0\varepsilon_0) \frac{\partial}{\partial t} (\mathbf{v} \times \mathbf{E}).\end{aligned}$$

EM energy in fluid

$$\partial_t u_{em} = -\nabla \cdot \mathbf{S} - \sigma \mathbf{E} \cdot (\mathbf{E} + \mu_0 \mathbf{v} \times \mathbf{H}) - \mathbf{v} \cdot \frac{\partial \mathbf{S}}{\partial t} - \mathbf{E} \cdot \mathbf{J}^{ext},$$

(20)

EM energy density

$$(21) \quad u_{em} = \frac{1}{2} (\epsilon \|\mathbf{E}\|^2 + \mu_0 \|\mathbf{H}\|^2) + 2\mathbf{v} \cdot \mathbf{S},$$

Poynting vector

$$(22) \quad \mathbf{S} = \mathbf{E} \times \mathbf{H}.$$

Time averaging

For periodic sources, the fields are periodic. We have

$$(23) \quad E_l(\mathbf{x}, t) = \operatorname{Re} \left\{ \hat{E}_l(\mathbf{x}) e^{i\omega t} \right\},$$

angular time frequency ω .

Time averaging:

$$\langle f(t) \rangle = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(t) dt.$$

Time averaging

For fields:

$$\langle \mathbf{E} \rangle = \langle \mathbf{H} \rangle = 0,$$

but:

$$\langle f_l \rangle = \partial_j \left[\mu_0 \mathbf{Re} \left\{ \hat{H}_l \bar{H}_j \right\} + \varepsilon_0 \mathbf{Re} \left\{ \hat{E}_l \bar{E}_j \right\} - \frac{1}{2} \delta_{lj} \left(\mu_0 |\hat{H}_l|^2 - \varepsilon_0 |\hat{E}_l|^2 \right) \right].$$

EM simulations: Meep

- Open source program, developed at MIT;
- Time domain simulations;
- Ability for frequency domain;
- Relatively easy C++ interface, easy to couple with other programs.
- Uses dimensionless Maxwell equations.

Maxwell equations: Dimensionless

Classical EM has four basic units:

- Electric current: I_0 in ampere,
- Distance: a in meter,
- Velocity: c in meter per second,
- Permittivity: ε in farad per meter.

Maxwell equations: Dimensionless

Dimensionless quantities:

$$\begin{aligned} E &= \frac{I_0}{a\epsilon c} E', & D &= \frac{I_0}{ac} D', \\ H &= \frac{I_0}{a} H', & B &= \frac{I_0}{ac^2\epsilon} B', \\ J &= \frac{I_0}{a^2} J', & K &= \frac{I_0}{a^2 c\epsilon} K', \\ \sigma &= \frac{\epsilon c}{a} \sigma', & \sigma_D &= \frac{c}{a} \sigma'_D, \end{aligned}$$

Furthermore:

$$x = ax', \quad t = \frac{a}{c}t'.$$

Maxwell equations: Dimensionless

Dimensionless equations:

$$\begin{aligned} -\epsilon_{ijk} \partial_{j'} H'_k + \partial_{t'} E'_i &= -J'_i{}^{\text{ind}} - J'_i{}^{\text{ext}}, \\ \epsilon_{ijk} \partial_{j'} E'_k + \partial_{t'} H'_i &= -K'_i{}^{\text{ind}} - K'_i{}^{\text{ext}}. \end{aligned}$$

Meep uses this form.

Navier Stokes equations

Conservation laws:

- Continuity equation: mass,
- Navier-Stokes: (linear) momentum,
- Energy equation: energy.

Continuity equation

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = 0,$$

- $\rho(\mathbf{x}, t)$ is the mass density,
- $\mathbf{v}(\mathbf{x}, t)$ is the fluid velocity.

Navier-Stokes equations

$$\rho (\partial_t v_i + v_j \partial_j v_i) = -\partial_i p + \partial_j T'_{ij} + f_i^{\mathbf{b}}.$$

- $p(\mathbf{x}, t)$ is the pressure,
- T'_{ij} is the deviatoric stress tensor,
- $f_i^{\mathbf{b}}$ are the body forces.

Stress tensor

For Newtonian fluids we can write:

$$T'_{ij} = 2\mu \left(e_{ij} - \frac{1}{3}\Delta\delta_{ij} \right),$$

where

- μ is the viscosity,
- $e_{ij} = \frac{1}{2} (\partial_j u_i + \partial_i u_j)$,
- $\Delta = e_{kk} = \partial_k v_k$

Incompressible

When flow is incompressible:

$$\nabla \cdot \mathbf{v} = 0.$$

Navier-Stokes simplifies:

$$\rho (\partial_t v_i + v_j \partial_j v_i) = -\partial_i p + \mu \partial_j^2 v_i + f_i^{\mathbf{b}}.$$

Energy equation

Conservation of energy:

$$\partial_t E + \partial_i (E v_i) = f_i^b v_i + \partial_j (T_{ij} v_i) + k \partial_i^2 T + q,$$

with

- E energy density,
- k thermal conductivity,
- q heat source density.

Energy density

For the energy density we have:

$$E = \rho \left(e + \frac{1}{2} v_i v_i \right),$$

with e the specific internal energy,

$$e = c_p T.$$

c_p is the specific heat (by constant pressure).

Free convection

If ΔT small,

$$T = T_0 + \Delta T, \quad \rho = \rho_0 + \Delta\rho.$$

Using linearisation we have:

$$\rho' = \left(\frac{\partial \rho_0}{\partial T} \right)_p T' = -\rho_0 \beta T',$$

So for the density:

$$\rho = \rho_0 (1 - \beta(T - T_0)).$$

Overview of equations

$$\partial_i v_i = 0,$$

$$\rho (\partial_t v_i + v_j \partial_j v_i) = -\partial_i p + \mu \partial_j^2 v_i + f_i^{\mathbf{b}},$$

$$\rho c_p (\partial_t T + v_i \partial_i T) = T_{ij} \partial_j v_i + k \partial_i^2 T + q,$$

$$\rho = \rho_0 (1 - \beta(T - T_0)).$$

Six equations, in six unknowns.

Solving the equations

TODO, first approach:

- Using OpenFoam for fluid simulations,
- Using Meep for EM simulations,
- Couple both programs.

Second approach:

Investigate if COMSOL can be of use.

Any suggestions?

Research questions

Once everything runs:

- Are the effects of the EM field significant?
- Which effect dominates, Lorentz force or heat convection?
- See if the model can be validated by experimental data.
- Investigate different materials.

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