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Abstract

In this paper, an eigenvalue analysis of the SIMPLE preconditioning for incompressible flow is presented. Some formulations have been set up to characterize the spectrum of the preconditioned matrix. This leads to a generalized eigenvalue problem. The generalized eigenvalue problem is investigated. Some eigenvalue bounds and the estimation for the spectral condition number in the symmetric case are given. Numerical tests are reported to illustrate the theoretical discussions.

Key words. preconditioning, SIMPLE preconditioner, eigenvalues, spectral analysis

AMS subject classification. 65N22

1   Introduction.

The steady state incompressible Navier-Stokes equations

$$\begin{cases} -\nu\Delta u + u \cdot \operatorname{grad} u + \operatorname{grad} p &= f, \\ -\operatorname{div} u &= 0, \end{cases}$$

combined with some boundary conditions, are widely used to simulate the incompressible flow of a fluid. Discretization and linearization of the equations leads to the following large sparse linear algebraic system

$$\begin{pmatrix} Q & G \\ G^T & O \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}, \tag{1.1}$$

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where $Q \in \mathbb{R}^{n \times n}, G \in \mathbb{R}^{n \times m}, m \leqslant n, \det(Q) \neq 0, \operatorname{rank}(G) = m$; $u \in \mathbb{R}^n$ and $p \in \mathbb{R}^m$ are the velocity vector and the pressure vector respectively. For problems with three space dimensions, iterative solvers are required. Preconditioning often determines the numerical performance of the Krylov subspace solvers [2].

In [9, 10], Vuik proposed GCR-SIMPLE(R) algorithm for solving the large linear system (1.1). The algorithm can be considered as a combination of the Krylov subspace method GCR [3] with the SIMPLE(R) algorithm[5]. In this combined algorithm, the SIMPLE(R) iteration is collaborated as a preconditioner with the GCR method. Numerical tests indicate that the SIMPLE(R) preconditioning is effective and competitive for practical use.

In this paper, we focus on the eigenvalue analysis of the SIMPLE preconditioned matrix $\widetilde{A}$. Two related formulations are derived to describe the spectrum of $\widetilde{A}$. The spectrum has some connection with that of the Schur complement of the matrix A. The relationship between the two different formulations has been investigated by using the theory of matrix singular value decomposition. A diagonal scaling technique proposed by Vuik[9] is studied. Some useful eigenvalue bounds have been got in symmetric situation. Numerical tests are used to illustrate the theoretical bounds.

In the remaining parts of this paper, the linear system (1.1) is abbreviated as $Ax = b$, where $A \in \mathbb{R}^{(n+m) \times (n+m)}, b \in \mathbb{R}^{n+m}$. Notations have the same meaning with references [10, 9]. $\sigma(A)$ represents the set of all eigenvalues of matrix A, for example. Besides, we assume that the matrix Q, its diagonal matrix $D := \operatorname{diag}(Q)$, and its Jacobi iteration matrix J ($J := D^{-1}(D - Q)$), are all nonsingular in this paper.

2 Formulations of the spectrum of the SIMPLE preconditioned matrix.

Consider the right preconditioning to the linear system (1.1)

$$AP^{-1}y = b, \quad x = P^{-1}y. \tag{2.1}$$

When the SIMPLE algorithm is used as preconditioning, it is equivalent to choose the preconditioner P^{-1} as[10, 11]

$$P^{-1} = BM^{-1} \quad , P = MB^{-1}, \tag{2.2}$$

where,

$$B = \begin{pmatrix} I & -D^{-1}G \\ O & I \end{pmatrix}, \quad M = \begin{pmatrix} Q & O \\ G^T & R \end{pmatrix}, \quad D = \operatorname{diag}(Q), \ R = -G^T D^{-1}G.$$

We call this preconditioning a SIMPLE preconditioning, and the preconditioner P^{-1} as SIMPLE preconditioner. For SIMPLE preconditioning, we have the following result:

Proposition 2.1. *If the right preconditioner P^{-1} is taken to be the matrix defined by (2.2), then the preconditioned matrix is*

$$\widetilde{A} := AP^{-1} = \begin{pmatrix} I - (I - QD^{-1})GR^{-1}G^TQ^{-1} & (I - QD^{-1})GR^{-1} \\ O & I \end{pmatrix}. \tag{2.3}$$

And, therefore, the spectrum of the SIMPLE *preconditioned matrix $\widetilde{A}$ is*

$$\sigma(\widetilde{A}) = \{1\} \cup \sigma(I - (I - QD^{-1})GR^{-1}G^TQ^{-1}). \tag{2.4}$$

Proof. It is easy to verify that

$$M^{-1} = \begin{pmatrix} Q^{-1} & O \\ -R^{-1}G^TQ^{-1} & R^{-1} \end{pmatrix}, \tag{2.5}$$

and

$$\begin{aligned} \widetilde{A} &= AP^{-1} = ABM^{-1} \\ &= \begin{pmatrix} Q & G \\ G^T & O \end{pmatrix} \begin{pmatrix} I & -D^{-1}G \\ O & I \end{pmatrix} \begin{pmatrix} Q^{-1} & O \\ -R^{-1}G^TQ^{-1} & R^{-1} \end{pmatrix} \\ &= \begin{pmatrix} I - (I - QD^{-1})GR^{-1}G^TQ^{-1} & (I - QD^{-1})GR^{-1} \\ O & I \end{pmatrix}. \end{aligned}$$

So, the fact about the spectrum of $\widetilde{A}$, described by (2.4), follows. $\qquad\square$

Now, we study the spectrum defined by (2.4) in more detail. By multiplying with matrices Q^{-1} and Q from the left- and right-hand side of the matrix $I - (I - QD^{-1})GR^{-1}G^TQ^{-1}$ respectively, we get

$$\begin{aligned} \sigma(I - (I - QD^{-1})GR^{-1}G^TQ^{-1}) &= \sigma(I - (Q^{-1} - D^{-1})GR^{-1}G^T) \\ &= \sigma(I - D^{-1}(D - Q)Q^{-1}GR^{-1}G^T), \\ &= \sigma(I - J\,Q^{-1}GR^{-1}G^T) \end{aligned}$$

in which, the matrix $J := D^{-1}(D - Q)$ is the *Jacobi* iteration matrix for the matrix Q. This observation leads to the following proposition:

Proposition 2.2. *For the* SIMPLE *preconditioned matrix $\widetilde{A}$,*

1. *1 is an eigenvalue with multiplicity at least of m, and*

2. *the remaining eigenvalues are $1 - \mu_i$, $i = 1, 2, \cdots, n$, where μ_i is the $i-$ th eigenvalue of the generalized eigenvalue problem*

$$Ex = \mu Zx, \tag{2.6}$$

 where,

$$E = GR^{-1}G^T \in \mathbb{R}^{n \times n}, \quad Z = QJ^{-1} \in \mathbb{R}^{n \times n}.$$

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Next, to investigate the spectrum of $\widetilde{A}$ more accurately, we derive another formulation of it. Consider the eigenvalue problem

$$\widetilde{A}x = \lambda x,$$

i.e.,

$$AP^{-1}x = \lambda x. \tag{2.7}$$

We know that AP^{-1} has the same spectrum as $P^{-1}A$ except for some possible zero eigenvalues [1]. When matrices A and P are both nonsingular, it holds that $\sigma(AP^{-1}) = \sigma(P^{-1}A)$. So, the eigenvalue problem (2.7) is equivalent to the generalized eigenvalue problem

$$Ax = \lambda Px. \tag{2.8}$$

Here,

$$A = \begin{pmatrix} Q & G \\ G^T & O \end{pmatrix},$$

and

$$P = MB^{-1} = \begin{pmatrix} Q & O \\ G^T & R \end{pmatrix} \begin{pmatrix} I & D^{-1}G \\ O & I \end{pmatrix} = \begin{pmatrix} Q & QD^{-1}G \\ G^T & O \end{pmatrix}.$$

The generalized eigenvalue problem (2.8) can be written as

$$\begin{pmatrix} Q & G \\ G^T & O \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \lambda \begin{pmatrix} Q & QD^{-1}G \\ G^T & O \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix}, \tag{2.9}$$

that is

$$\begin{cases} Qu + Gp & = \lambda(Qu + QD^{-1}Gp), \\ G^T u & = \lambda G^T u. \end{cases}$$

Multiply by Q^{-1} from the left to the first equation, and re-arrange the two equations as

$$\begin{cases} (1 - \lambda)u & = (\lambda D^{-1} - Q^{-1})Gp, \\ G^T(1 - \lambda)u & = 0. \end{cases} \tag{2.10}$$

From (2.10), we see that 1 is an eigenvalue of (2.9).Note that the matrix $D^{-1} - Q^{-1} = -JQ^{-1}$ is nonsingular by assumption, from the right-hand side of the first equation of (2.10)$(\lambda = 1)$,we can see that the eigenvectors corresponding to eigenvalue 1 are:

$$v_i = \begin{pmatrix} u_i \\ 0 \end{pmatrix} \in \mathbb{R}^{(n+m)}, u_i \in \mathbb{R}^n, \ i = 1, 2, \cdots, n,$$

where, $\{u_i\}_{i=1}^n$ is a basis of $\mathbb{R}^n$.

For $\lambda \neq 1$, it follows from the second equation in (2.10) that $G^T u = 0$. Multiplying the first equation in (2.10) with G^T shows that

$$0 = -G^T Q^{-1} Gp + \lambda G^T D^{-1} Gp,$$
$$-G^T Q^{-1} Gp = -\lambda G^T D^{-1} Gp.$$

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It follows that
$$Sp = \lambda Rp,$$
in which, $S = -G^T Q^{-1} G \in \mathbb{R}^{m \times m}$ is the *Schur* complement of the matrix A, and $R = -G^T D^{-1} G \in \mathbb{R}^{m \times m}$.

To conclude the above analysis, the following proposition is derived.

Proposition 2.3. *For the* SIMPLE *preconditioned matrix* $\widetilde{A}$,

1. *1 is an eigenvalue with multiplicity of n, and*

2. *the remaining eigenvalues are defined by the generalized eigenvalue problem*

$$Sp = \lambda Rp. \tag{2.11}$$

In the following section, we investigate the generalized eigenvalue problems in more detail.

3 Further investigation on the spectrum of $\widetilde{A}$

In section 2, two different generalized eigenvalue problems (2.6) and (2.11) have been derived to describe the spectrum of $\widetilde{A}$. In this section, we shall show that the two generalized eigenvalue problems are closely related.

Firstly, we investigate the generalized eigenvalue problem (2.11). Re-write matrix R as

$$R = -G^T D^{-1} G = -(D^{-\frac{1}{2}} G)^T (D^{-\frac{1}{2}} G).$$

Making the singular value decomposition of the matrix $D^{-\frac{1}{2}} G \in \mathbb{R}^{n \times m}$, we have

$$D^{-\frac{1}{2}} G = U \Sigma V^T, \tag{3.1}$$

in which, $U \in \mathbb{R}^{n \times n}, V \in \mathbb{R}^{m \times m}$ are unitary matrices, i.e., $U^T U = I \in \mathbb{R}^{n \times n}, \quad V^T V = I \in \mathbb{R}^{m \times m}$, and

$$\Sigma = \begin{pmatrix} \sigma_1 & & & \\ & \sigma_2 & & O \\ & & \ddots & \\ & & & \sigma_m \\ & O & & \end{pmatrix} \in \mathbb{R}^{n \times m},$$

$\sigma_i, i = 1, 2, \cdots, m$, are the singular values of the matrix $D^{-\frac{1}{2}} G$, which are all positive numbers since $\text{rank}(D^{-\frac{1}{2}} G) = m$. So,

$$\begin{aligned}
G &= D^{\frac{1}{2}} U \Sigma V^T, \\
R &= -(U \Sigma V^T)^T (U \Sigma V^T) = -V \Sigma^T \Sigma V^T, \\
S &= -G^T Q^{-1} G \\
&= -(D^{\frac{1}{2}} U \Sigma V^T)^T Q^{-1} (D^{\frac{1}{2}} U \Sigma V^T) \\
&= -V \Sigma^T U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U \Sigma V^T.
\end{aligned}$$

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It follows that
$$R^{-1}S = V(\Sigma^T\Sigma)^{-1}\Sigma^T U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U \Sigma V^T. \tag{3.2}$$

To study the generalized eigenvalue problem (2.6), by using the same singular value decomposition for matrix $D^{-\frac{1}{2}}G$, we have

$$\begin{aligned}
E\ &= GR^{-1}G^T \\
&= (D^{\frac{1}{2}}U\Sigma V^T)(-V(\Sigma^T\Sigma)^{-1}V^T)(D^{\frac{1}{2}}U\Sigma V^T)^T \\
&= -D^{\frac{1}{2}}U\Sigma(\Sigma^T\Sigma)^{-1}\Sigma^T U^T D^{\frac{1}{2}}.
\end{aligned}$$

The matrix Z is a notation for matrix QJ^{-1}, so

$$Z^{-1} = JQ^{-1} = D^{-1}(D-Q)Q^{-1} = (Q^{-1} - D^{-1}).$$

Finally, we get
$$Z^{-1}E = -(Q^{-1} - D^{-1})D^{\frac{1}{2}}U\Sigma(\Sigma^T\Sigma)^{-1}\Sigma^T U^T D^{\frac{1}{2}}. \tag{3.3}$$

Multiplying by $U^T D^{\frac{1}{2}}$ and $D^{-\frac{1}{2}}U$ to (3.3) from the left-side and right-side respectively, a spectrum equivalent matrix is produced as

$$U^T D^{\frac{1}{2}} Z^{-1} E D^{-\frac{1}{2}} U = -U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U \Sigma(\Sigma^T\Sigma)^{-1}\Sigma^T + \Sigma(\Sigma^T\Sigma)^{-1}\Sigma^T.$$

We denote this equation by
$$U^T D^{\frac{1}{2}} Z^{-1} E D^{-\frac{1}{2}} U = -MN + N, \tag{3.4}$$

in which,
$$M = U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U \in \mathbb{R}^{n \times n},$$

and

$$\begin{aligned}
N\ &= \Sigma(\Sigma^T\Sigma)^{-1}\Sigma^T \\
&= \begin{pmatrix} \sigma_1 & & & & \\ & \sigma_2 & & & \\ & & \ddots & & \\ & & & \sigma_m & \\ & & O & & O \end{pmatrix}_{n \times m} \begin{pmatrix} \frac{1}{\sigma_1^2} & & & \\ & \frac{1}{\sigma_2^2} & & \\ & & \ddots & \\ & & & \frac{1}{\sigma_m^2} \end{pmatrix}_{m \times m} \begin{pmatrix} \sigma_1 & & & & \\ & \sigma_2 & & & \\ & & \ddots & & O \\ & & & \sigma_m & \end{pmatrix}_{m \times n} \\
&= \begin{pmatrix} 1 & & & & \\ & 1 & & O & \\ & & \ddots & & \\ & & & 1 & \\ & O & & & O \end{pmatrix}_{n \times n} = \begin{pmatrix} I_m & O \\ O & O \end{pmatrix} \in \mathbb{R}^{n \times n}.
\end{aligned}$$

Partitioning matrix M according to the structure of N, (3.4) can be written in a sub-matrix form

$$
\begin{aligned}
U^T D^{\frac{1}{2}} Z^{-1} E D^{-\frac{1}{2}} U \ &= -MN + N \\
&= -\begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} I_m & O \\ O & O \end{pmatrix} + \begin{pmatrix} I_m & O \\ O & O \end{pmatrix} \\
&= \begin{pmatrix} I_m - M_{11} & O \\ -M_{21} & O \end{pmatrix}.
\end{aligned}
\tag{3.5}
$$

Its characteristic polynomial is

$$
\det(\mu I - U^T D^{\frac{1}{2}} Z^{-1} E D^{-\frac{1}{2}} U) = \mu^{n-m} \det((\mu - 1)I_m + M_{11}).
$$

So, we get to know that 0 is an eigenvalue of $Z^{-1}E$ with multiplicity of $n - m$, and the remaining eigenvalues are $\mu_i = 1 - \eta_i, i = 1, 2, \cdots, m$, where η_i is the i−th nonzero eigenvalue of the sub-matrix M_{11}. From (3.5),η_i is also an eigenvalue of MN at the same time, since that

$$
\det(\eta I - MN) = \eta^{n-m} \det(\eta I_m - M_{11}).
$$

By Proposition 2.2, we have

$$
\sigma(\widetilde{A}) = \{1\} \cup \{1 - \mu_i\} = \{1\} \cup \{\eta_i\},
\tag{3.6}
$$

in which, the eigenvalue 1 has the multiplicity of $m + (n - m) = n$, and $\eta_i \in \sigma(MN), \eta_i \neq 0, i = 1, 2, \cdots, m$.

On the other hand, if we denote

$$
T_1 := U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U\Sigma \in \mathbb{R}^{n \times m},
$$

and

$$
T_2 := (\Sigma^T \Sigma)^{-1} \Sigma^T \in \mathbb{R}^{m \times n},
$$

then $MN = T_1 T_2$. We know that $T_1 T_2 \in \mathbb{R}^{n \times n}$ and $T_2 T_1 \in \mathbb{R}^{m \times m}$ have the same spectrum except for the possible zero eigenvalue [1, pp.69]. The spectrum of $T_2 T_1$ is

$$
\begin{aligned}
\sigma(T_2 T_1) \ &= \sigma((\Sigma^T \Sigma)^{-1} \Sigma^T U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U\Sigma) \\
&= \sigma(V(\Sigma^T \Sigma)^{-1} \Sigma^T U^T D^{\frac{1}{2}} Q^{-1} D^{\frac{1}{2}} U\Sigma V^T) \\
&= \sigma(R^{-1}S).
\end{aligned}
$$

The last equation is based on the fact of equation (3.2). This relation motivates the following proposition.

Proposition 3.1. *For the two generalized eigenvalue problem (2.6) and (2.11), suppose that $\mu_i \in \sigma(Z^{-1}E), i = 1, 2, \cdots, n$, and $\lambda_i \in \sigma(R^{-1}S), i = 1, 2, \cdots, m$, the relationship between the two problems is that $\mu = 0$ is an eigenvalue of (2.6) with multiplicity of $n - m$, which can be denoted as $\mu_{m+1} = \mu_{m+2} = \cdots = \mu_n = 0$, and that $\lambda_i = 1 - \mu_i, i = 1, 2, \cdots, m$, holds for the remaining m eigenvalues.*

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4 The influence of the diagonal scaling.

Vuik [9, 10] proposed a diagonal scaling strategy for practical implementation of the SIMPLE preconditioning. Scale the coefficient matrix A by (left) multiplying the diagonal matrix

$$\widehat{D} := \begin{pmatrix} D^{-1} & O \\ O & D_R^{-1} \end{pmatrix}, \tag{4.1}$$

where,

$$D = \mathrm{diag}(Q), D_R = \mathrm{diag}(R), R = -G^T D^{-1} G.$$

After the scaling, the coefficient matrix becomes to be

$$\mathcal{A} := \widehat{D}A = \begin{pmatrix} D^{-1}Q & D^{-1}G \\ D_R^{-1}G^T & O \end{pmatrix}. \tag{4.2}$$

At this moment,

$$\mathcal{D} = \mathrm{diag}(D^{-1}Q) = I \in \mathbb{R}^{n \times n}, \mathcal{R} = -(D_R^{-1}G^T)\mathcal{D}^{-1}(D^{-1}G) = D_R^{-1}R \in \mathbb{R}^{m \times m},$$

and

$$\mathcal{B} = \begin{pmatrix} I & -D^{-1}G \\ O & I \end{pmatrix}, \mathcal{M} = \begin{pmatrix} D^{-1}Q & O \\ D_R^{-1}G^T & \mathcal{R} \end{pmatrix}, \mathcal{M}^{-1} = \begin{pmatrix} Q^{-1}D & O \\ -\mathcal{R}^{-1}D_R^{-1}G^T Q^{-1}D & \mathcal{R}^{-1} \end{pmatrix}.$$

The SIMPLE preconditioned matrix now is

$$\begin{aligned} \widetilde{\mathcal{A}} &= \mathcal{A}\mathcal{B}\mathcal{M}^{-1} \\ &= \begin{pmatrix} D^{-1}Q & D^{-1}G \\ D_R^{-1}G^T & O \end{pmatrix} \begin{pmatrix} I & -D^{-1}G \\ O & I \end{pmatrix} \begin{pmatrix} Q^{-1}D & O \\ -\mathcal{R}^{-1}D_R^{-1}G^T Q^{-1}D & \mathcal{R}^{-1} \end{pmatrix} \\ &= \begin{pmatrix} \widetilde{\mathcal{A}}_{11} & \widetilde{\mathcal{A}}_{12} \\ \widetilde{\mathcal{A}}_{21} & \widetilde{\mathcal{A}}_{22} \end{pmatrix}, \end{aligned}$$

in which, by doing some elementary matrix calculation, these sub-matrices are:

$$\begin{aligned} \widetilde{\mathcal{A}}_{11} &= I + D^{-1}\big[QD^{-1}G\mathcal{R}^{-1}D_R^{-1}G^T Q^{-1} - G\mathcal{R}^{-1}D_R^{-1}G^T Q^{-1}\big]D \\ &= I - D^{-1}Q(Q^{-1} - D^{-1})GR^{-1}G^T Q^{-1}D, \\ \widetilde{\mathcal{A}}_{12} &= -D^{-1}QD^{-1}G\mathcal{R}^{-1} + D^{-1}G\mathcal{R}^{-1} = D^{-1}(I - QD^{-1})GR^{-1}, \\ \widetilde{\mathcal{A}}_{21} &= D_R^{-1}G^T Q^{-1}D + D_R^{-1}G^T D^{-1}G\mathcal{R}^{-1}D_R^{-1}G^T Q^{-1}D = O, \\ \widetilde{\mathcal{A}}_{22} &= -D_R^{-1}G^T D^{-1}G\mathcal{R}^{-1} = I. \end{aligned}$$

Finally, it follows that

$$\widetilde{\mathcal{A}} = \begin{pmatrix} I - D^{-1}Q(Q^{-1} - D^{-1})GR^{-1}G^T Q^{-1}D & D^{-1}(I - QD^{-1})GR^{-1}D_R \\ O & I \end{pmatrix}. \tag{4.3}$$

Comparing the matrix $\widetilde{\mathcal{A}}$ in (4.3) with the matrix $\widetilde{A}$ defined by (2.3), we find that the spectra of both matrices are exactly the same. So, theoretically speaking, there is no influence to the spectrum of the SIMPLE preconditioned matrix by the diagonal scaling (4.1).

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5   Some eigenvalue bounds for symmetric case.

In this section, we assume that Q is symmetric positive definite, which corresponds to the cases when term $u \operatorname{grad} u$ is deleted from Navier-Stokes equations in incompressible flow. In this case, the coefficient matrix A is symmetric and indefinite.

Consider the generalized eigenvalue problem (2.11)

$$Sp = \lambda Rp. \tag{5.1}$$

It is obvious that the problem $-Sp = -\lambda Rp$ is completely equivalent to the problem $Sp = \lambda Rp$. Since both $-S$ and $-R$ are s.p.d. matrices, we call (5.1) as a s.p.d. generalized eigenvalue problem by neglecting the negative signs in both sides. For the s.p.d. generalized eigenvalue problem, the extreme eigenvalues ($\lambda_{\max}$ and $\lambda_{\min}$) are the extreme values of [1, pp.379]:

$$\frac{p^T S p}{p^T R p} = \frac{p^T G^T Q^{-1} G p}{p^T G^T D^{-1} G p}, \quad p \neq 0, p \in \mathbb{R}^m, \tag{5.2}$$

which is the ratio of the Rayleigh quotients of S and R. So,

$$\lambda_{\max} = \max_{p \neq 0} \frac{p^T G^T Q^{-1} G p}{p^T G^T D^{-1} G p} = \max_{p \neq 0} \frac{(Gp)^T Q^{-1} (Gp)}{(Gp)^T D^{-1} (Gp)}. \tag{5.3}$$

Since that the matrix G has column full rank, i.e. $\operatorname{rank}(G) = m$, $Gp = 0$ if and only if $p = 0$. Denoting $y = Gp$, it follows that

$$\lambda_{\max} \leq \max_{y \neq 0} \frac{y^T Q^{-1} y}{y^T D^{-1} y}. \tag{5.4}$$

Let μ_1, μ_n be the largest and the smallest eigenvalues of the matrix Q, and d_1, d_n be the largest and the smallest diagonal elements of Q respectively, then

$$\lambda_{\max} \leqslant \frac{d_1}{\mu_n}. \tag{5.5}$$

It is easy to show that

$$\lambda_{\min} \geqslant \frac{d_n}{\mu_1} \tag{5.6}$$

by a similar argument.

So, combining (5.5), (5.6) and Proposition 2.3, we get the following bounds for the eigenvalues of the preconditioned matrix $\widetilde{A}$:

$$\min\left\{1, \frac{d_n}{\mu_1}\right\} \leqslant \lambda \leqslant \max\left\{1, \frac{d_1}{\mu_n}\right\}, \quad \forall \lambda \in \sigma(\widetilde{A}). \tag{5.7}$$

If the both sides of (5.7) are taken to be $\frac{d_n}{\mu_1}$ and $\frac{d_1}{\mu_n}$ respectively, then

$$\kappa(\widetilde{A}) = \frac{\lambda_{\max}}{\lambda_{\min}} \leqslant \frac{d_1}{d_n} \cdot \frac{\mu_1}{\mu_n} = \frac{d_1}{d_n} \kappa(Q), \tag{5.8}$$

where, $\kappa(\cdot)$ represents for the (spectral) condition number.

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6 Numerical examples.

Two numerical test results are reported here to illustrate the discussions above.

Example 6.1. *In this example, the coefficient matrix is taken from a discretised Navier-Stokes equations on a 16×16 grid [10](length= 2, $\nu = 1$). The dimensions are $n = 544, m = 256$, and $n + m = 800$. $A \in \mathbb{R}^{800 \times 800}$ is a nonsymmetric matrix.*

The eigenvalues of the preconditioned matrix $\widetilde{A}$ were computed by both Proposition 2.2 and Proposition 2.3. The computing results were the same, which coincided with the theoretical analysis. Spectra of A and $\widetilde{A}$ are plotted in Figure 6.1, and some extreme eigenvalues are listed in Table 6.1.
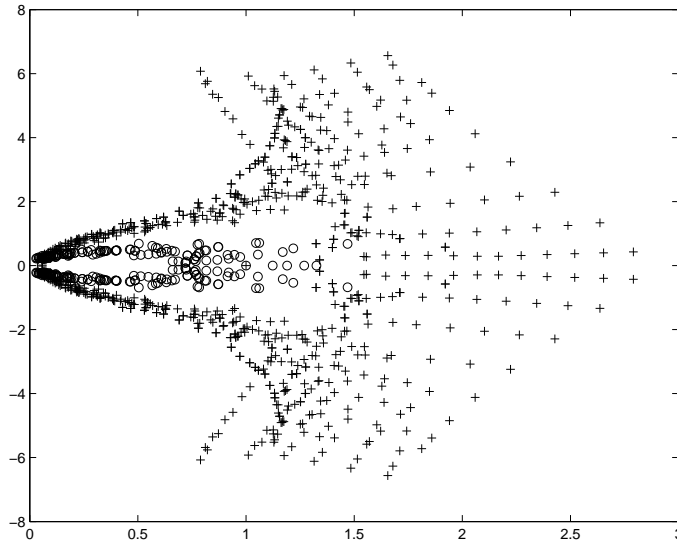

Figure 6.1. Spectrum of A and $\widetilde{A}$.
The '+' represents for the eigenvalues of A, while 'o' for that of the preconditioned $\widetilde{A}$.

Table 6.1. The extreme eigenvalues of A and $\widetilde{A}$.

matrix	$\max \Re(\lambda_i)$	$\min \Re(\lambda_i)$	$\max \Im(\lambda_i)$	$\max |\lambda_i|$	$\min |\lambda_i|$
A	2.79074	0.03559	6.56341	6.76892	0.06018
$\widetilde{A}$	1.46960	0.03000	0.70700	1.61894	0.21395

Example 6.2. *The matrix A is obtained from a discretised Stokes equation on a 16×16 grid by removing the Dirichlet boundary conditions. The resulted coefficient matrix $A \in \mathbb{R}^{800 \times 800}$ is symmetric, and $Q \in \mathbb{R}^{544 \times 544}$ is a s.p.d. matrix.*

The extreme eigenvalues of A and $\widetilde{A}$ are listed in Table 6.2.

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Table 6.2. The extreme eigenvalues of A and $\widetilde{A}$ for example 6.2.

matrix	$\lambda_{\min}$	$\min|\lambda_i|$	$\lambda_{\max}$	$\kappa(\cdot)$
A	-23.4555	0.0501	25.3762	1.7295×10^3
$\widetilde{A}$	0.5049	0.5049	46.7880	344.1452
Q	0.0154	0.0154	2.5477	232.9809
$\mathrm{diag}(Q)$	0.9600	0.9600	1.6000	1.6667

From example 6.1, we can see that the eigenvalues of the SIMPLE preconditioned matrix $\widetilde{A}$ are clustered in a smaller region in the right -half plane. The results of example 6.2 agree with the theoretical eigenvalue bounds in section 5, which are:

$$\frac{0.96}{2.547} = 0.377, \text{ and } \frac{1.6}{0.0154} = 103.9.$$

Both examples indicate that the spectrum could be effectively improved by using the SIMPLE preconditioner.

7 Concluding remarks and future work.

We have derived two formulations to describe the spectrum of the SIMPLE preconditioned matrix $\widetilde{A}$. These theoretical results could be helpful to achieve new insights for this pre-conditioning. The methodology in this paper is instructive for the eigenvalue analysis for this type of preconditioning (for example, the SIMPLER preconditioning). The eigenvalue bounds in the symmetric case could be useful for evaluating the efficiency of the SIMPLE preconditioned iterative solvers for Stokes equations.

The results for general non-symmetric matrix in this paper mainly have some theoretical meaning. More accurate and more practical estimations about the spectrum of $\widetilde{A}$ need to be done. The main issues towards this aim are the investigations to the specific generalized eigenvalue problems (2.6) and (2.11). Pseudo-spectra analysis [7, 8] might be needed to analyze these non-symmetric problems.

SIMPLER preconditioning seems to be more effective [10]. Eigenvalue analysis for this preconditioning is also necessary.

Elman, Silvester, and Wathen [4, 6] proposed a kind of preconditioners based on some approximations to the Schur complement of A. The comparisons of the theoretical proper-ties and the numerical performance of this kind of preconditioners with that of the SIMPLE type preconditioners should be done.

References

[1] O. Axelsson, Iterative Solution Methods, Cambridge University Press, Cambridge, UK, 1994.

[2] A. M. Bruaset, A Survey of Preconditioned Iterative Methods, Longman Scientific & Technical Press, UK, 1995.

[3] S. C. Eisenstat, H. C. Elman, and H. C. Schultz, Variational iterative methods for non-symmetric systems of linear equations, SIAM J. Numer. Anal., 20(1983), pp.345-357.

[4] H. C. Elman, D. J. Silvester, and A. J. Wathen, Performance and Analysis of Saddle Point Preconditioners for the Discrete Steady-State Navier-Stokes Equations, Numer. Math., 90(2002), pp.665-688.

[5] S. V. Patankar, Numerical Heat Transfer and Fluid Flow, McGraw-Hill, New York, 1980.

[6] D. J. Silvester, H. C. Elman, D. Kay and A. J. Wathen, Efficient preconditioning of the linearized Navier-Stokes equations, J. Comput. Appl. Math., 128 (2001), no. 1-2, pp.261-279.

[7] L. N. Trefethen, Pseudospectra of linear operators, SIAM Review 39 (1997), pp.383-406.

[8] L. N. Trefethen, Spectra and pseudospectra, In M. Ainsworth, J. Levesley, and M. Marletta, eds., The Graduate Student's Guide to Numerical Analysis '98, Springer-Verlag, Berlin, 1999, pp. 217-250.

[9] C. Vuik, A. Saghir, and G. P. Boerstoel, The Krylov accelerated SIMPLE(R) method for flow problems in industrial furnaces, Inter. J. Numer. Methods in Fluids, 33(2000), pp.1027-1040.

[10] C. Vuik and A. Saghir, The Krylov accelerated SIMPLE(R) method for incompressible flow, TUD Report No. 02-01, 2002.

[11] P. Wesseling, An Introduction to Multigrid Methods, JohnWiley & Sons, New York, 1991.