

A scalable parallel solver for Helmholtz problems

Cornelis Vuik, Jinqiang Chen, Vandana Dwarka

Technische Universiteit Delft

Aim and Impact

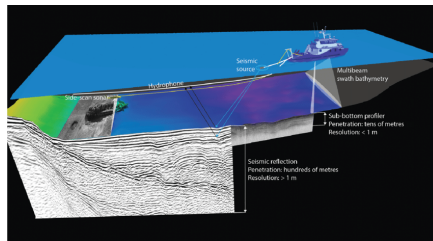
- **Contribute** to broad research on parallel scalable iterative solvers for time-harmonic wave problems
- This presentation: **matrix-free parallelization**
 - Complex shift Laplace Preconditioner (CSLP)
 - Deflation methods
 - Parallel performance

Introduction - the Helmholtz Problem

 The Helmholtz equation (describing time-harmonic waves) + BCs

$$-\Delta u(\mathbf{x}) - k(\mathbf{x})^2 u(\mathbf{x}) = g(\mathbf{x}), \text{ on } \Omega \subseteq \mathbb{R}^n$$

- > Δ - Laplace operator, $u(\mathbf{x})$ - Fourier-space representation of the wave function
- > $k(\mathbf{x})$ - **wavenumber**, $k(\mathbf{x}) = (2\pi f)/c(\mathbf{x})$, where f - **frequency**, c - wave velocity
- > Applications in **seismic exploration**, medical imaging, antenna synthesis, etc.



 Larisa, High-performance implementation of Helmholtz equation with absorbing boundary conditions.

<http://www.math.chalmers.se/~larisa/www/MasterProjects/HelmholtzABSbc.pdf>

 M. Jakobsson, et al (2016). Mapping submarine glacial landforms using acoustic methods. Geological Society.

Introduction - Challenges

 Linear system from discretization

$$Au = b$$

➤ A is real, sparse, symmetric, normal, and **indefinite**; **non-Hermitian** with Sommerfeld BCs

? Direct solver or iterative solver

⚠ **Accuracy and pollution error** ($k^3 h^2 < 1$): finer grid (3D) $\Rightarrow$ larger linear system

🔧 Memory-efficient methods; High-Performance Computing (HPC)

⚠ **Negative & positive eigenvalues**: larger wavenumber $\Rightarrow$ more iterations

🔧 Preconditioner: Complex Shifted Laplace Preconditioner (CSLP)

🔧 (Higher-order) Deflation

⚠ **Parallelism**

Aim

💡 A wavenumber-independent convergent and parallel scalable solver

Introduction - Metrics

➤ Convergence metric:

- Krylov-based solvers, GMRES-type: the number of iterations (#iter); IDR(s): the number of matrix-vector multiplications (#Matvec)

➤ Scalability:

- Strong scaling: the number of processors is increased while the problem size remains constant
- Weak scaling: the problem size increases along with the number of tasks, so the computation per task remains constant
- Wall-clock time: t_w ; number of processors: np
- Speedup: $S_p = \frac{t_{w,r}}{t_{w,p}}$, $E_P = \frac{S_p}{np/np_r} = \frac{t_{w,r} \cdot np_r}{t_{w,p} \cdot np}$

Introduction - Numerical Models

- Model problems on a rectangular domain Ω with boundary $\Gamma = \partial\Omega$

$$-\Delta u(\mathbf{x}) - k(\mathbf{x})^2 u(\mathbf{x}) = \delta(\mathbf{x} - \mathbf{x}_0), \text{ on } \Omega$$

$$\frac{\partial u(\mathbf{x})}{\partial \vec{n}} - ik(\mathbf{x})u(\mathbf{x}) = 0, \text{ on } \Gamma$$

- Constant wavenumber: $k(\mathbf{x}) = k$
- Non-constant wavenumber: Wedge, Marmousi problem
- Finite-difference discretization on a uniform grid with grid size h . (2D example)
- Laplace operator:

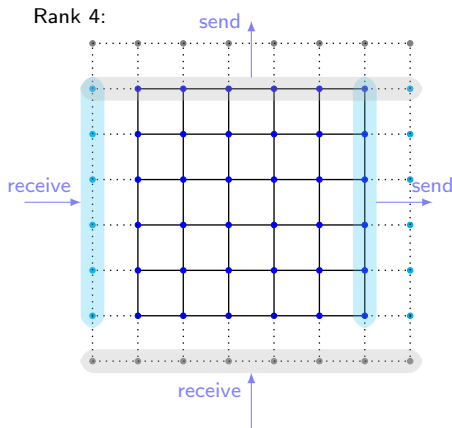
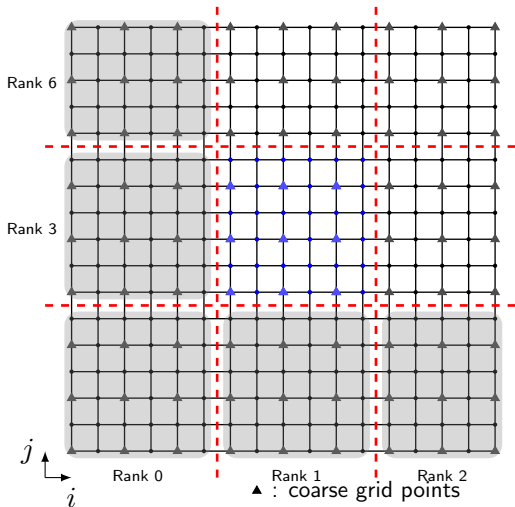
$$-\Delta_h \mathbf{u} \approx \frac{-u_{i,j-1} - u_{i-1,j} + 4u_{i,j} - u_{i+1,j} - u_{i,j+1}}{h^2}$$

- Sommerfeld BCs: a ghost point

$$\frac{\partial u}{\partial \vec{n}}(0, y_j) - ik(0, y_j)u(0, y_j) \approx \frac{u_{0,j} - u_{2,j}}{2h} - ik_{1,j}u_{1,j} = 0 \Rightarrow u_{0,j} = u_{2,j} + 2hik_{1,j}u_{1,j}$$

Framework - Distributed data structure

- Vector $\mathbf{u} \Leftarrow$ 2D array: $\mathbf{u}(1:N_x, 1:N_y) \Leftarrow$ each sub-domain: $\mathbf{u}(1-LAP:nx+LAP, 1-LAP:ny+LAP)$
- Operations (e.g. matvec, dot-product, vector update) perform on each $\mathbf{u}(1:nx, 1:ny)$ simultaneously



Framework - Matrix-free operations

- ▶ Perform computations with a matrix without explicitly forming or storing the matrix
⇒ Reduce memory requirements

Matrix-vector multiplication

If a matrix can be represented by a so-called stencil notation

$$[A] = \begin{bmatrix} a_{-1,1} & a_{0,1} & a_{1,1} \\ a_{-1,0} & a_{0,0} & a_{1,0} \\ a_{-1,-1} & a_{0,-1} & a_{1,-1} \end{bmatrix},$$

Then $\mathbf{v} = A\mathbf{u}$ can be computed by

$$v_{i,j} = \sum_{p=-1}^1 \sum_{q=-1}^1 a_{p,q} u_{i+p,j+q}$$

with the help of a ghost point on the physical boundary and one overlapping grid point.

Framework - Matrix-free operations

i Stencil notation

> Laplace operator:

$$[-\Delta_h] = \frac{1}{h^2} \begin{bmatrix} 0 & -1 & 0 \\ -1 & 4 & -1 \\ 0 & -1 & 0 \end{bmatrix}$$

> “Wavenumber operator”:

$$[\mathcal{I}_h \mathbf{k}^2] = \begin{bmatrix} 0 & 0 & 0 \\ 0 & k_{i,j}^2 & 0 \\ 0 & 0 & 0 \end{bmatrix} \stackrel{const}{=} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} k^2$$

> $A\mathbf{u} = \mathbf{b}$:

$$[A_h] = [-\Delta_h] - [\mathcal{I}_h \mathbf{k}^2]$$

- **Speed up** convergence of Krylov subspace methods by **Preconditioning**
- Solve $M^{-1}Au = M^{-1}b$
- Complex Shifted Laplace Preconditioner (CSLP)

$$M_h = -\Delta_h - (\beta_1 - \beta_2 i) \mathcal{I}_h \mathbf{k}^2, \quad (\beta_1, \beta_2) \in [0, 1], \quad \text{e.g. } \beta_1 = 1, \beta_2 = 0.5$$

☑ Stencil notation

- Solve $Mx = u$ by multigrid method (V-cycle) $\Rightarrow x \approx M^{-1}u$
 - **Vertex-centered** coarsening based on the **global** grid
 - Damped Jacobi smoother (easy to parallelize)
 - Full-weight restriction I_h^{2h} & linear interpolation I_{2h}^h

$$[I_h^{2h}] = \frac{1}{16} \begin{bmatrix} 1 & 2 & 1 \\ 2 & 4 & 2 \\ 1 & 2 & 1 \end{bmatrix}_h^{2h}, \quad [I_{2h}^h] = \frac{1}{4} \begin{bmatrix} 1 & 2 & 1 \\ 2 & 4 & 2 \\ 1 & 2 & 1 \end{bmatrix}_{2h}^h$$

- Coarse-grid operator obtained by **re-discretization**
 - ☑ Stencil notation: $[M_{2h}]$ similar to $[M_h]$

CSLP - Cons

- Increasing $k \Rightarrow$ eigenvalues move fast towards **origin**
- Too many iterations for high frequency
- **Project** unwanted eigenvalues to **zero** $\Rightarrow$ **Deflation**

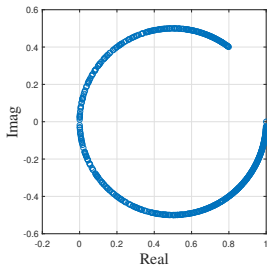
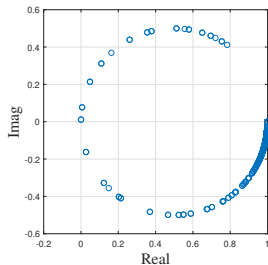


Figure: $\sigma \left(M_{(1,0.5)}^{-1} A \right)$ for $k = 20$ (left) and $k = 80$ (right)

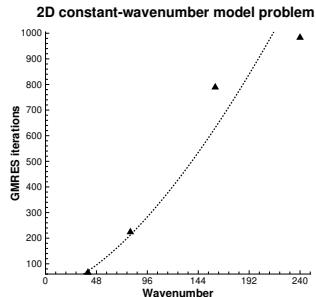


Figure: #Iter increases with k

Deflation - introduction

- **Project** unwanted eigenvalues to **zero** $\Rightarrow$ **Deflation**

- Deflation preconditioning: solve $PA\hat{u} = Pb$

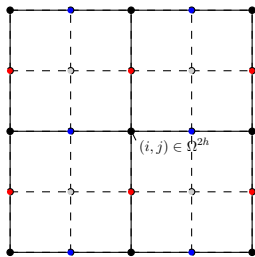
$$P = I - AQ, \quad \text{where } Q = ZE^{-1}Z^T, \quad E = Z^T AZ$$
$$A \in \mathbb{R}^{n \times n}, Z \in \mathbb{R}^{m \times n}$$

- Columns of Z span deflation subspace
- Ideally Z contains **eigenvectors**
- In practice **approximations**: inter-grid vectors from **multigrid**
- Adapted Deflation Variant 1 (A-DEF1): $P_{A-DEF1} = M_{(\beta_1, \beta_2)}^{-1}P + Q$
 - > Combined with the standard preconditioner CSLP
- Linear approximation basis deflation vectors $\rightarrow$ **higher-order** deflation vectors (Adapted Preconditioned DEF, **APD**)
 - > wavenumber-independent convergence

Higher-order deflation vectors

- 2D: the higher-order interpolation & restriction has 5×5 stencil
 - **Two overlapping** grid points are needed

$$[Z] = \frac{1}{64} \begin{bmatrix} 1 & 4 & 6 & 4 & 1 \\ 4 & 16 & 24 & 16 & 4 \\ 6 & 24 & 36 & 24 & 6 \\ 4 & 16 & 24 & 16 & 4 \\ 1 & 4 & 6 & 4 & 1 \end{bmatrix} \begin{matrix} h \\ \\ \\ 2h \end{matrix}, \quad [Z^T] = \frac{1}{64} \begin{bmatrix} 1 & 4 & 6 & 4 & 1 \\ 4 & 16 & 24 & 16 & 4 \\ 6 & 24 & 36 & 24 & 6 \\ 4 & 16 & 24 & 16 & 4 \\ 1 & 4 & 6 & 4 & 1 \end{bmatrix} \begin{matrix} 2h \\ \\ \\ h \end{matrix}$$



- • • • : fine grid points $\in \Omega^h$
- : coarse grid points $\in \Omega^{2h}$

Figure: The allocation map of interpolation operator

Matrix-free two-level deflation

$$P = I - AQ, \quad \text{where } Q = ZE^{-1}Z^T, \quad E = Z^T AZ$$

- With matrix constructed, $E = Z^T AZ$, so-called Galerkin Coarsening

Matrix-free coarse grid operation $y = Ex$?

- Straightforward Galerkin Coarsening operator;

$$x_1 = Zx, \quad x_2 = A_h x_1, \quad y = Z^T x_2 \Rightarrow y = Ex$$

- unacceptable **computational cost** for consideration of multilevel method

- Re-discretization:

- 💡 **ReD-02**: The same as the fine grid

- 💡 **ReD-04**: Fourth-order re-discretization of the Laplace operator

$$[E] = \frac{1}{12 \cdot (2h)^2} \begin{bmatrix} 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & -16 & 0 & 0 \\ 1 & -16 & 60 & -16 & 1 \\ 0 & 0 & -16 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \end{bmatrix} - \mathcal{I}_{2h} \mathbf{k}_{2h}^2$$

Matrix-free two-level deflation

💡 **ReD-Glk**: Re-discretized scheme (stencil) from the result of Galerkin coarsening

$$[-\Delta_{2h}] = \frac{1}{(2h)^2} \cdot \frac{1}{256} \begin{bmatrix} -3 & -44 & -98 & -44 & -3 \\ -44 & -112 & 56 & -112 & -44 \\ -98 & 56 & 980 & 56 & -98 \\ -44 & -112 & 56 & -112 & -44 \\ -3 & -44 & -98 & -44 & -3 \end{bmatrix}$$

$$\Rightarrow -\Delta_{2h} u_{2h} = -4 \frac{\partial^2 u}{\partial x^2} - 4 \frac{\partial^2 u}{\partial y^2} - \left(\frac{13}{48} \frac{\partial^4 u}{\partial x^4} + \frac{1}{2} \frac{\partial^4 u}{\partial x^2 \partial y^2} + \frac{13}{48} \frac{\partial^4 u}{\partial y^4} \right) (\mathbf{2h})^2 + \mathcal{O}(h^4)$$

$$[\mathcal{I}_{2h} \mathbf{k}_{2h}^2] = \frac{1}{64^2} \begin{bmatrix} 1 & 28 & 70 & 28 & 1 \\ 28 & 784 & 1960 & 784 & 28 \\ 70 & 1960 & 4900 & 1960 & 70 \\ 28 & 784 & 1960 & 784 & 28 \\ 1 & 28 & 70 & 28 & 1 \end{bmatrix} \mathbf{k}_{2h}^2$$

$$\Rightarrow [E] = [-\Delta_{2h}] - [\mathcal{I}_{2h} \mathbf{k}_{2h}^2]$$

❓ Boundary conditions - 💡 **ReD-O2** on the boundary grid points

Two-level deflation - overall algorithm

➤ **Flexible** GMRES-type methods → allow for variable preconditioner

Algorithm Two-level deflation **FGMRES**

- 1: Choose u_0 and dimension k of the Krylov subspace.
 - 2: Define $(k+1) \times k \bar{H}_k$ and initialize to zero
 - 3: Compute $r_0 = b - Au_0$, $\beta = \|r_0\|$, $v_1 = r_0/\beta$;
 - 4: **for** $j = 1, 2, \dots, k$ or until convergence **do**
 - 5: $\hat{v}_j = Z^T v_j$ ▷ Precondition starts
 - 6: $\tilde{v} \approx E^{-1} \hat{v}$ ▷ Solved by GMRES approximately, preconditioned by CSLP, **tol**= 10^{-1}
 - 7: $t = Z \tilde{v}$
 - 8: $s = At$
 - 9: $\tilde{r} = v_j - s$
 - 10: $r \approx M^{-1} \tilde{r}$ ▷ Approximated by one multigrid V-cycle
 - 11: $x_j = r + t$ ▷ Precondition ends
 - 12: $w = Ax_j$
 - 13: **for** $i := 1, 2, \dots, j$ **do**
 - 14: $h_{i,j} = (w, v_i)$
 - 15: $w := w - h_{i,j} v_i$
 - 16: **end for**
 - 17: $h_{j+1,j} := \|w\|_2$, $v_{j+1} = w/h_{j+1,j}$; $X_k = [x_1, \dots, x_k]$; $\bar{H}_k = \{h_{i,j}\}_{1 \leq i \leq j+1, 1 \leq j \leq m}$
 - 18: **end for**
 - 19: $u_k = u_0 + X_k y_k$ where $y_k = \arg \min_y \|\beta e_1 - \bar{H}_k y\|$
-

Tolerance for the coarse grid problem

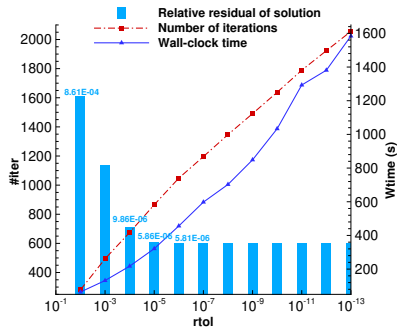


Figure: APD preconditioned GMRES

- Outer #iter keeps **constant**
- 10^{-6} for the coarse grid problem is **essential**

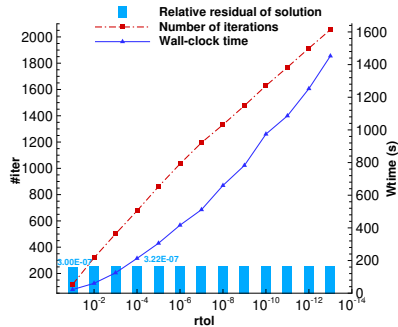


Figure: APD preconditioned FGMRES

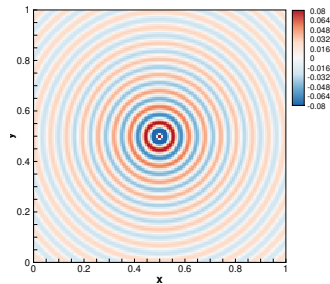
- Outer #iter keeps 10, but 11 for 10^{-1}
- **One more** outer iteration, **much less** Wtime

Convergence - Constant wavenumber

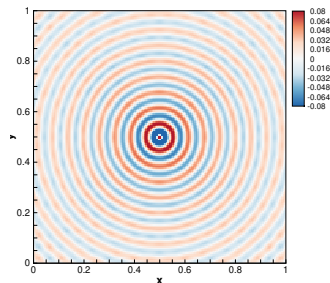
Table: The number of iterations required by using APD-GMRES.

Grid size	k	kh	ReD- $\mathcal{O}2$	ReD- $\mathcal{O}4$	ReD-Glk
65×65	40	0.625	20	17	9
129×129	80	0.625	30	18	9
257×257	160	0.625	87	19	9
513×513	320	0.625	319	23	10
1025×1025	640	0.625	1099	34	11
2049×2049	1280	0.625	3417	79	13
129×129	40	0.3125	18	18	7
257×257	80	0.3125	19	18	7
513×513	160	0.3125	21	18	7
1025×1025	320	0.3125	28	20	6
2049×2049	640	0.3125	53	23	6

- ☑ $Ex = Z^T A_h Z x$: #iter=**7** for $kh = 0.625$, **5** for $kh = 0.3125$
- ☑ ReD- $\mathcal{O}4$ better than ReD- $\mathcal{O}2$
- ☑ ReD-Glk: close to **wavenumber independence**



(a) Exact solution



(b) $kh = 0.625$

Convergence - 2D Wedge

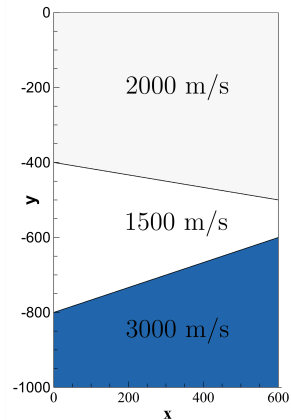


Figure: Wedge problem

Convergence - 2D Wedge

Table: The number of iterations required by using APD-GMRES.

Grid size	f	kh	ReD- $\mathcal{O}2$	ReD- $\mathcal{O}4$	ReD-Glk
73×121	10	0.35	22	22	9
145×241	20	0.35	28	27	9
289×481	40	0.35	31	29	9
577×961	80	0.35	37	30	9
1153×1921	160	0.35	58	34	8

- ✔ $Ex = Z^T A_h Zx$: #iter=**6**
- ✔ ReD- $\mathcal{O}4$ better than ReD- $\mathcal{O}2$
- ✔ ReD-Glk: **wavenumber independence** although it is derived from constant wavenumber

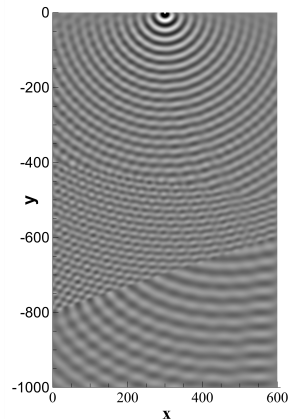
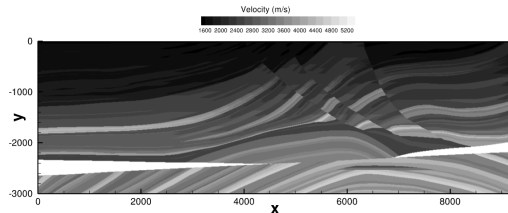
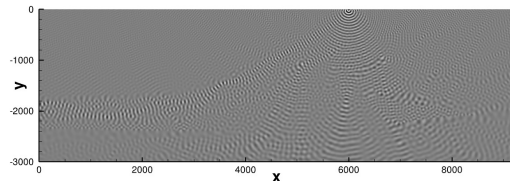


Figure: Waves pattern at 80 Hz

Convergence - Marmousi



(a) Marmousi problem



(b) Wave pattern at $f = 40$ Hz

Table: The number of iterations required by using APD-FGMRES.

Grid size	f	kh	ReD- $\mathcal{O}2$	ReD- $\mathcal{O}4$	ReD-Glk
737×241	10	0.5236	40	33	11
1473×481	20	0.5236	71	35	12
2945×961	40	0.5236	233	41	12

- ✓ $Ex = Z^T A_h Zx$: #iter=8
- ✓ ReD-Glk: close to **wavenumber independence** for **highly heterogeneous** media
- ✓ Many iterations are required to solve the coarse grid problem (**in parentheses**)
 $\Rightarrow$ Apply two-level method **recursively** $\Rightarrow$ Multilevel deflation

Parallel performance - Weak scaling

- › APD using ReD-Glk
- › DelftBlue, GNU Fortran 8.5.0, Open MPI 4.1.1

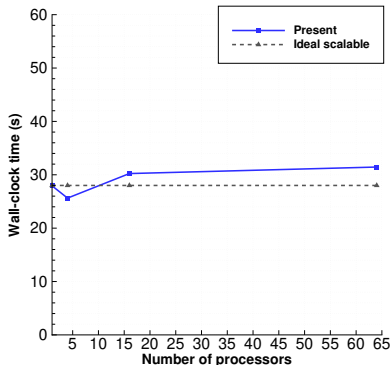


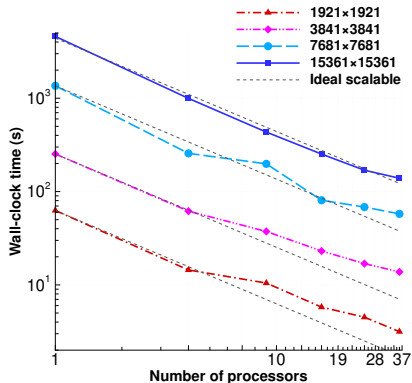
Figure: Weak scaling for constant-wavenumber problem with $k = 160$ and a grid size of 1024×1024 per processes.

Table: Weak scaling for model problems with non-constant wavenumber.

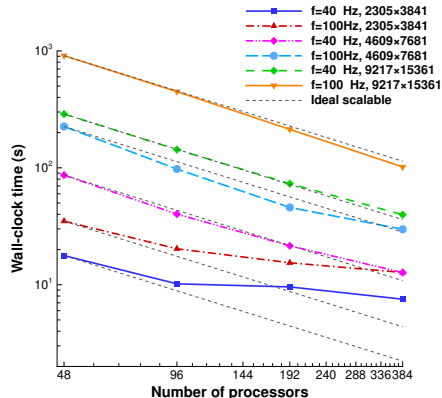
grid size	np	#iter	CPU time (s)
Wedge, $f = 40$ Hz			
577×961	6	10 (46)	4.86
1153×1921	24	10 (43)	5.75
Marmousi, $f = 10$ Hz			
737×241	3	11 (63)	10.55
1473×481	12	10 (58)	12.08
2945×961	48	10 (58)	17.72

👍 Close to weak scalability

Parallel performance - Strong scaling



(a) Constant-wavenumber problem with $k = 200$



(b) Wedge problem with $f = 40$ Hz and $f = 100$ Hz

☑ Good strong scalability for large problems (larger computation/communication ratio)

Multilevel deflation

Table: The number of outer iterations required to solve the Wedge problem with $kh = 0.17$ by using the **multilevel (six-level) APD-FGMRES**.

Grid size	f (Hz)	Outer #iter (L2 #iter)
289×481	20	11 (3)
577×961	40	12 (4)
1153×1921	80	12 (7)
2305×3841	160	13 (13)
4609×7681	320	14 (27)
9217×15361	640	17 (47)

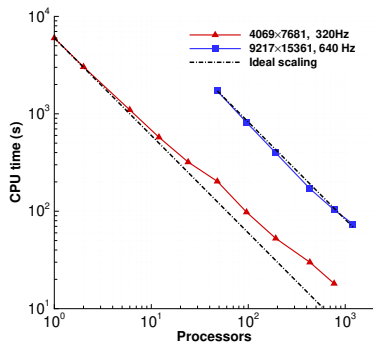


Figure: Strong scaling for Wedge problem

- ✓ The number of iterations weakly depends on the frequency
- ✓ Good performance for massively parallel computing

Conclusions

- ✓ Parallelization of higher-order deflation preconditioned Krylov solvers
- ✓ Matrix-free implementation with wavenumber-independent convergence
- ✓ Parallel framework with good weak and strong scalability

Further reading:

- 📄 Dwarka, V., Vuik, C.: Scalable convergence using two-level deflation preconditioning for the Helmholtz equation, *SIAM Journal on Scientific Computing*, 42(2020), A901-A928.
- 📄 Dwarka, V., Vuik, C.: Scalable multi-level deflation preconditioning for highly indefinite time-harmonic waves, *Journal of Computational Physics*, 469(2022), 111327.
- 📄 Chen, J., Dwarka, V., Vuik, C.: A matrix-free parallel solution method for the three-dimensional heterogeneous Helmholtz equation, *Electronic Transactions on Numerical Analysis*, 59 (2023), 270–294.
- 📄 Chen, J., Dwarka, V., Vuik, C.: A matrix-free parallel two-level deflation preconditioner for two-dimensional heterogeneous Helmholtz problems, *Journal of Computational Physics* 514 (2024): 113264.
- 📄 Chen, J., Dwarka, V., Vuik, C.: Matrix-Free Parallel Scalable Multilevel Deflation Preconditioning for Heterogeneous Time-Harmonic Wave Problems, *Journal of Scientific Computing*, 102.2 (2025), 1-33.

Thanks!