Multi-level Krylov: the next generation Helmholtz solver

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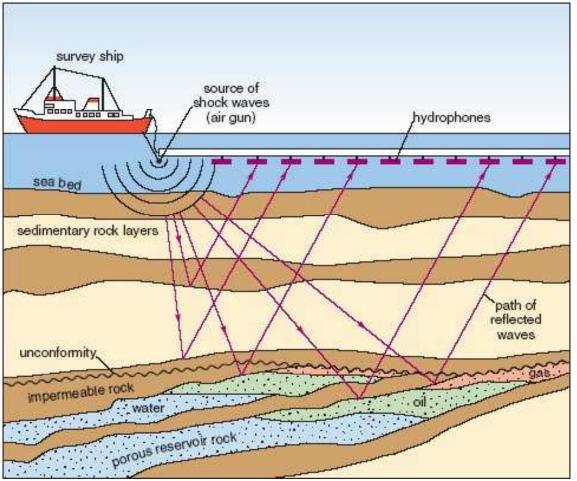
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Marine Seismic

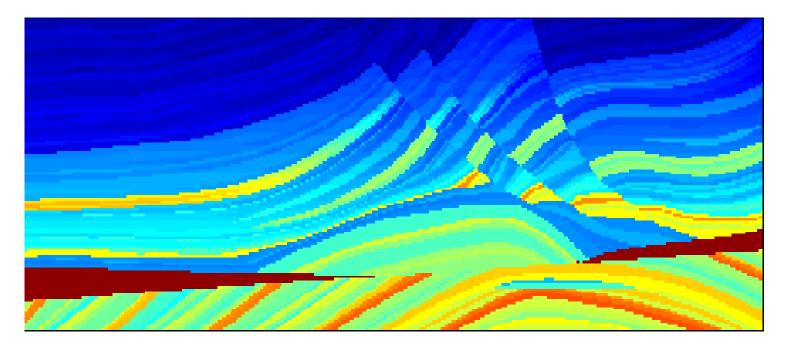


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hard Marmousi Model

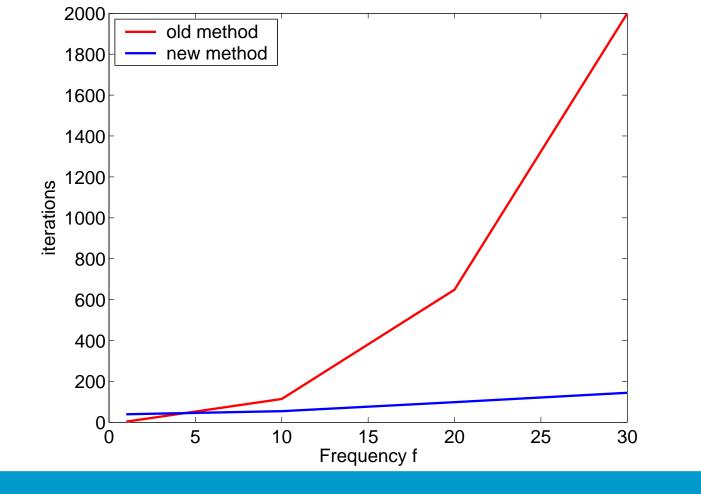


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hard Marmousi Model (2006)





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1. Introduction

The Helmholtz equation without damping

 $-\Delta \mathbf{u}(x,y) - k^2(x,y)\mathbf{u}(x,y) = \mathbf{g}(x,y) \text{ in } \Omega$

 $\mathbf{u}(x,y)$ is the pressure field,

 $\mathbf{k}(x,y)$ is the wave number,

 $\mathbf{g}(x,y)$ is the point source function and

 Ω is the domain. Absorbing boundary conditions are used on $\Gamma.$

$$\frac{\partial \mathbf{u}}{\partial n} - \iota \mathbf{u} = 0$$

n is the unit normal vector pointing outwards on the boundary.

Perfectly Matched Layer (PML) and Absorbing Boundary Layer (ABL)

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Problem description

• Second order Finite Difference stencil:

$$\begin{array}{ccc}
-1 \\
-1 & 4 - k^2 h^2 & -1 \\
-1 & -1
\end{array}$$

- Linear system Au = g: properties
 Sparse & complex valued
 Symmetric & Indefinite for large k
- For high resolution a very fine grid is required: 10 20 gridpoints per wavelength $\rightarrow A$ is extremely large!
- Is traditionally solved by a Krylov subspace method, which exploits the sparsity.

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2. Preconditioning

Equivalent linear system $M_1^{-1}AM_2^{-1}\tilde{x} = \tilde{b}$, where $M = M_1 \cdot M_2$ is the preconditioning matrix and

$$\tilde{x} = M_2 x, \quad \tilde{b} = M_1 b.$$

Requirements for a preconditioner

- better spectral properties of $M^{-1}A$
- cheap to perform $M^{-1}r$.

Spectrum of A is $\{\mu_i - k^2\}$, with k a given constant and μ_i are the

eigenvalues of the Laplace operator. Note that $\mu_1 - k^2$ may be negative.

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Preconditioning (overview)

ILU Meijerink and van der Vorst, 1977

ILU(tol) Saad, 2003

SPAI Grote and Huckle, 1997

Multigrid Lahaye, 2001 Elman, Ernst and O' Leary, 2001

AILU Gander and Nataf, 2001 analytic parabolic factorization

ILU-SV Plessix and Mulder, 2003 separation of variables

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Preconditioning (Laplace type)

Laplace operatorBayliss and Turkel, 1983Definite HelmholtzLaird, 2000Shifted LaplaceY.A. Erlangga, C. Vuik and C.W.Oosterlee, 2003

Shifted Laplace preconditioner (SLP)

$$M \equiv -\Delta - (\beta_1 - \mathbf{i}\beta_2)k^2, \ \beta_1, \beta_2 \in \mathbb{R}.$$

Condition $\beta_1 \leq 0$ is used to ensure that M is a (semi) definite operator.

- $\rightarrow \beta_1, \beta_2 = 0$: Bayliss and Turkel
- $ightarrow eta_1 = -1, eta_2 = 0$: Laird
- $\rightarrow \beta_1 = 1, \beta_2 = 0.5$: Y.A. Erlangga, C. Vuik and C.W.Oosterlee

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3. Numerical experiments

Example with constant k in Ω

Iterative solver: Bi-CGSTAB

Preconditioner: Shifted-Laplace operator, discretized using the same method as the Helmholtz operator.

k	ILU(0.01)	M_0	M_{-1}	M_i
5	9	13	13	13
10	25	29	28	22
15	47	114	45	26
20	82	354	85	34
30	139	> 1000	150	52

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Spectrum of SLP

References: Manteuffel, Parter, 1990; Yserentant, 1988

Since $L \equiv -\Delta$ is SPD we have the following eigenpairs

$$Lv_j = \lambda_j v_j$$
, where, $\lambda_j \in \mathbb{R}^+$

The eigenvalues σ_i of the preconditioned matrix satisfy

$$(L-z_1I)v_j = \sigma_j(L-z_2I)v_j.$$

Theorem 1

Provided that $z_2 \neq \lambda_j$, the relation

$$\sigma_j = rac{\lambda_j - z_1}{\lambda_j - z_2}$$
 holds.

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Spectrum of SLP

Theorem 2

If $\beta_2 = 0$, the eigenvalues $\sigma_r + i\sigma_i$ are located on the straight line in the complex plane given by

$$\beta_1 \sigma_r - (\alpha_1 - \alpha_2) \sigma_i = \beta_1.$$

Theorem 3

If $\beta_2 \neq 0$, the eigenvalues $\sigma_r + i\sigma_i$ are on the circle in the complex plane with center *c* and radius *R*:

$$c = \frac{z_1 - \bar{z}_2}{z_2 - \bar{z}_2}, \quad R = \left| \frac{z_2 - z_1}{z_2 - \bar{z}_2} \right|.$$

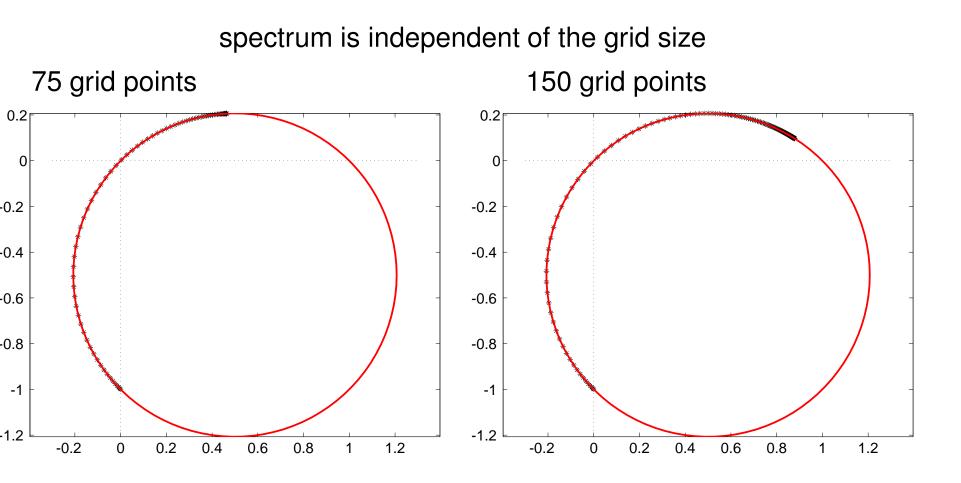
Note that if $\beta_1\beta_2 > 0$ the origin is not enclosed in the circle.

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Eigenvalues for Complex preco k = 100



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Inner iteration

Possible solvers for solution of Mz = r:

- ILU approximation of *M*
- inner iteration with ILU as preconditioner
- Multigrid

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Multigrid components

- geometric multigrid
- Gauss-Seidel with red-black ordering
- matrix dependent interpolation, full weighting restriction
- Galerkin coarse grid approximation



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Numerical results for a wedge problem

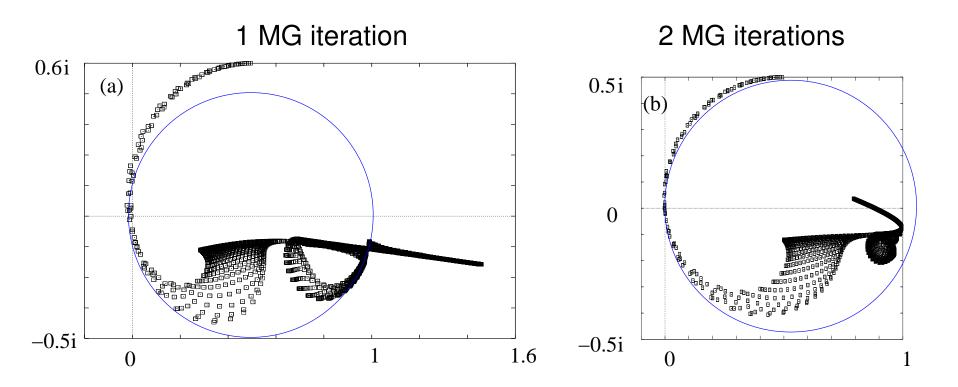
k_2	10	20	40	50	100
grid	32^{2}	64^{2}	128^{2}	192^{2}	384^{2}
No-Prec	201(0.56)	1028(12)	5170(316)	_	—
ILU(A,0)	55(0.36)	348(9)	1484(131)	2344(498)	—
ILU(A, 1)	26(0.14)	126(4)	577(62)	894(207)	—
ILU(M,0)	57(0.29)	213(8)	1289(122)	2072(451)	—
ILU(M, 1)	28(0.28)	116(4)	443(48)	763(191)	2021(1875)
MG(V(1,1))	13(0.21)	38(3)	94(28)	115(82)	252(850)

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K

Spectrum with inner iteration



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4. Second Level Precond. (2008-2014)

Summary so far

- ILU and variants
- From Laplace to complex Shifted Laplace Preconditioner (2005)
- Shifted Laplace Preconditioner (SLP)

$$M := -\Delta \mathbf{u} - (\beta_1 - \iota \beta_2) k^2 \mathbf{u}$$

- Results show: $(\beta_1, \beta_2) = (1, 0.5)$ is the shift of choice
- Properties of SLP?



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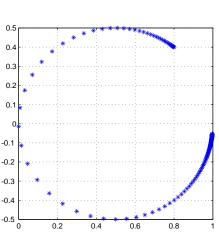
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Shifted Laplace Preconditioner (SLP)

- Introduces damping, Multi-grid approximation is possible
- The modulus of all eigenvalues of the preconditioned operator is bounded by 1
- Small eigenvalues move to zero, as k increases.

Spectrum of $M^{-1}(1, 0.5)A$ for

k = 30 and



k = 120

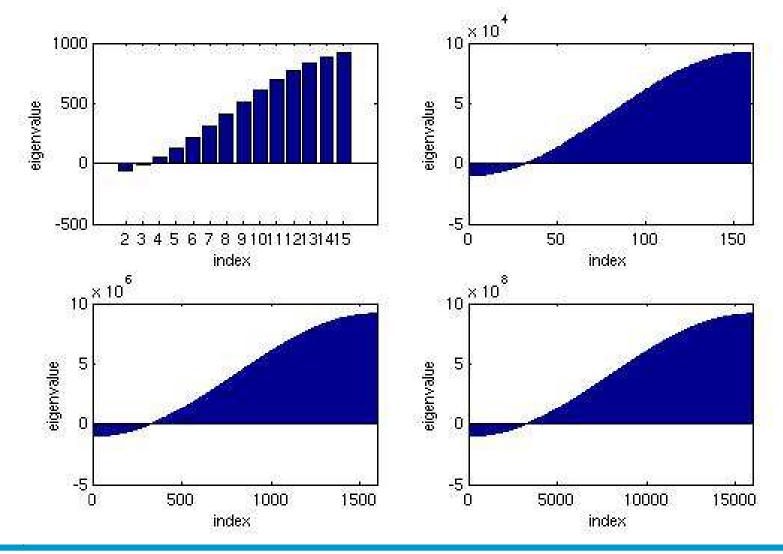


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Spectrum as function of k





Deflation: or two-grid method

Deflation, a projection preconditioner

P = I - AQ, with $Q = ZE^{-1}Z^T$ and $E = Z^TAZ$

where,

 $Z \in \mathbb{R}^{n \times r}$, with deflation vectors $Z = [z_1, ..., z_r]$, $rank(Z) = r \le n$

Along with a traditional preconditioner M, deflated preconditioned system reads

 $PM^{-1}Au = PM^{-1}g.$

Deflation vectors shifted the eigenvalues to zero.



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Deflation for Helmholtz

With choice of multigrid inter-grid transfer operator (Prolongation) as deflation matrix, i.e. $Z = I_h^{2h}$ and $Z^T = I_{2h}^h$ then

 $P_h = I_h - A_h Q_h$, with $Q_h = I_h^{2h} A_{2h}^{-1} I_{2h}^h$ and $A_{2h} = I_{2h}^h A_h I_h^{2h}$

where

- P_h can be interpreted as a coarse grid correction and
- Q_h as the coarse grid operator





Deflation: ADEF1

Deflation can be implemented combined with SLP M_h ,

 $M_h^{-1}P_hA_hu_h = M_h^{-1}P_hg_h$

 $A_h u_h = g_h$ is preconditioned by the two-level preconditioner $M_h^{-1} P_h$.

For large problems, A_{2h} is too large to invert exactly. Inversion of A_{2h} is sensitive, since P_h deflates the spectrum to zero.

To do: Solve A_{2h} iteratively to a required accuracy on certain levels, and shift the deflated spectrum to λ_h^{max} by adding a shift in the two level preconditioner. This leads to the **ADEF1** preconditioner

 $P_{(h,ADEF1)} = M_h^{-1} P_h + \lambda_h^{max} Q_h$

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Deflation: MLKM

Multi Level Krylov Method ^{*a*}, take $\hat{A}_h = M_h^{-1}A_h$, and define \hat{P}_h by using \hat{A}_h (instead of A_h) will be

$$\hat{P}_h = I_h - \hat{A}_h \hat{Q}_h,$$

where

$$\hat{Q}_h = I_h^{2h} \hat{A}_{2h}^{-1} I_{2h}^h$$
 and $\hat{A}_{2h} = I_{2h}^h \hat{A}_h I_h^{2h} = I_{2h}^h (M_h^{-1} A_h) I_h^{2h}$

Construction of coarse matrix A_{2h} at level 2h costs inversion of preconditioner at level h. Approximate A_{2h}

Ideal	Practical
$\hat{A}_{2h} = I_{2h}^h (M_h^{-1} A_h) I_h^{2h}$	$\hat{A}_{2h} = I_{2h}^h (M_h^{-1} A_h) I_h^{2h}$
	$\hat{A}_{2h} \approx I_{2h}^{h} I_{h}^{2h} M_{2h}^{-1} A_{2h}$

^aErlangga, Y.A and Nabben R., ETNA 2008

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5. Fourier Analysis of two-level methods

Dirichlet boundary conditions for analysis. With above deflation,

 $\operatorname{spec}(PM^{-1}A) = f(\beta_1, \beta_2, k, h)$

is a complex valued function.

Setting kh = 0.625,

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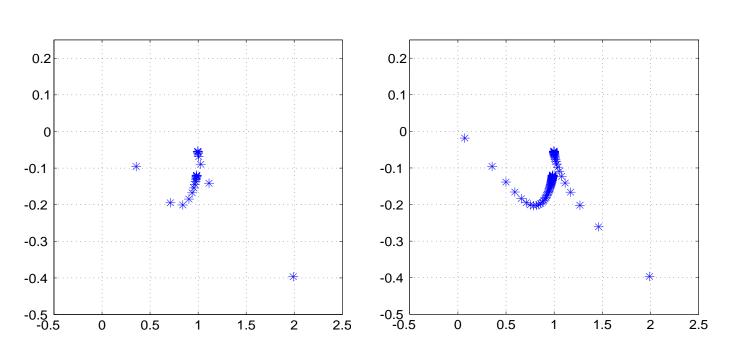
- Spectrum of $PM^{-1}A$ with shifts (1, 0.5) is clustered around 1 with a few outliers.
- Spectrum remains almost the same, when the imaginary shift for the preconditioner is varied from 0.5 to 1.



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Fourier Analysis

<u>ADEF1:</u> Analysis shows spectrum clustered around 1 with few outliers.



 $k = 30 \qquad \qquad k = 120$

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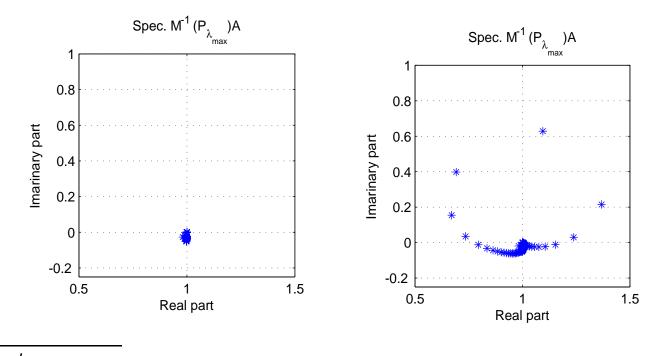
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Fourier Analysis

Spectrum of Helmholtz preconditioned by <u>MLKM</u> b , k = 160 and 20 gp/wl Ideal Practical



^bTwo-level

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6. Numerical results (no pollution)

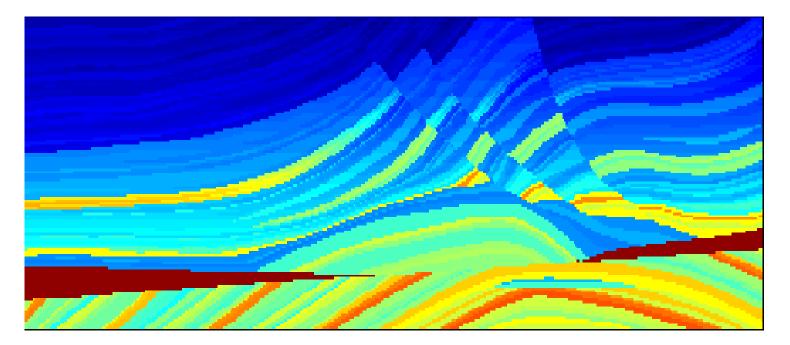
k	N ($k^3h^2 \le 0.625$)	iter	N ($kh \le 0.625$)	iter
10	44	4	16	4
20	116	4	32	5
40	320	4	64	5
80	800	4	128	6
100	1268	4	160	7
200	3572	4	320	8
400	10124	4	340	10
500	14144	4	800	12
800	28628	4	1280	15
1000	40004	4	1600	17

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hard Marmousi Model



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hard Marmousi Model, PETSc solver

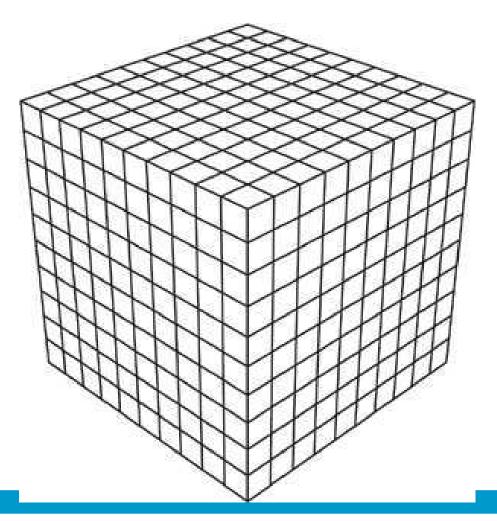
kh = 0.39, Bi-CGSTAB for SLP, FGMRES(20) for ADEF1(8,2,1)

Frequency f	Solve Time		Iterations	
	SLP-F ADEF1-F		SLP-F	ADEF1-F
1	1.22	5.07	13	7
10	10.18	9.43	112	13
20	72.16	60.32	189	22
40	550.20	426.79	354	39



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Cube with constant k



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Cube with constant k

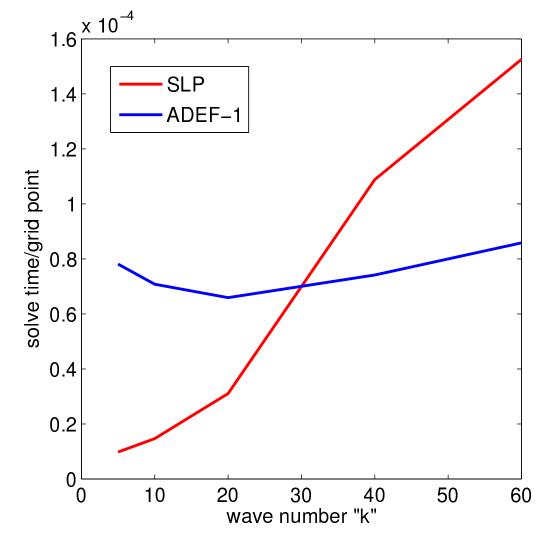
Wave number	Solve Time		Iterations	
k	SLP-F ADEF1-		SLP-F	ADEF1-F
5	0.04	0.32	7	8
10	0.48	2.32	9	9
20	8.14	17.28	20	9
40	228.29	155.52	70	10
60	1079.99	607.45	97	11

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Cube with constant k





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Cube with variable k

Grid size h is such that $kh \approx 0.625$

k	CLSP(time)	ADEF1(time)	CLSP	ADEF1
5	0.09	0.24	9	11
10	1.07	1.94	15	12
20	16.7	18.9	32	16
40	1304	214	331	24





Cube with variable k

Grid size h is such that $kh \approx 0.3125$

k	CLSP(time)	ADEF1(time)	CLSP	ADEF1
5	0.6	1.4	9	9
10	7.5	10.04	14	9
20	324.1	79.2	72	9
30	3810.9	361.7	285	11





7. Conclusions

- Without deflation, when imaginary shift is increased in SLP, spectrum remains bounded above 1, but lower part moves to zero.
- With deflation the convergence is nearly independent of the imaginary shift.
- With deflation the convergence is initially weakly depending on k.
 For large k is scales again linearly.
- With deflation the CPU time is less than without deflation.
- The convergence of ADEF1 and the practical variant of MLKM are similar.



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