## Deflation type methods combined with shifted Laplace preconditioners for the Helmholtz equation

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Kees Vuik, Abdul Sheikh and Domenico Lahaye
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## Application: geophysical survey

hard Marmousi Model


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## 1. Problem: The Helmholtz equation

The Helmholtz equation without damping

$$
-\Delta \mathbf{u}(x, y)-k^{2}(x, y) \mathbf{u}(x, y)=\mathbf{g}(x, y) \text { in } \Omega
$$

$\mathbf{u}(x, y)$ is the pressure field, $\mathbf{k}(x, y)$ is the wave number, $\mathrm{g}(x, y)$ is the point source function and
$\Omega$ is the domain. Absorbing boundary conditions are used on $\Gamma$.

$$
\frac{\partial \mathbf{u}}{\partial n}-\iota \mathbf{u}=0
$$

$n$ is the unit normal vector pointing outwards on the boundary.
Perfectly Matched Layer (PML) and Absorbing Boundary Layer (ABL)

## Problem description

- Second order Finite Difference stencil:

$$
\left[\begin{array}{ccc} 
& -1 & \\
-1 & 4-k^{2} h^{2} & -1 \\
& -1 &
\end{array}\right]
$$

- Linear system $A_{h} u_{h}=g_{h}$ : properties

Sparse \& complex valued
Symmetric \& Indefinite for large $k$

- For high resolution a very fine grid is required: $30-60$ gridpoints per wavelength (or $\approx 5-10 \times k$ ) $\rightarrow A_{h}$ is extremely large!
- Traditionally solved by a Krylov subspace method, which exploits the sparsity.


## 2. Shifted Laplace Preconditioner

Laplace operator Bayliss and Turkel, 1983
Definite Helmholtz Laird, 2000
Shifted Laplace Y.A. Erlangga, C. Vuik and C.W.Oosterlee, 2003
Shifted Laplace preconditioner

$$
M \equiv-\Delta-\left(\beta_{1}-\mathbf{i} \beta_{2}\right) k^{2}, \quad \beta_{1}, \beta_{2} \in \mathbb{R}
$$

If $\beta_{1} \leq 0$ holds than $M$ is a (semi) definite operator.
$\rightarrow \beta_{1}, \beta_{2}=0 \quad: \quad$ Bayliss and Turkel
$\rightarrow \beta_{1}=1, \beta_{2}=0 \quad: \quad$ Laird
$\rightarrow \beta_{1}=-1, \beta_{2}=0.5 \quad$ : Y.A. Erlangga, C. Vuik and C.W. Oosterlee

## Numerical results for a wedge problem

| $k_{2}$ | 10 | 20 | 40 | 50 | 100 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| grid | $32^{2}$ | $64^{2}$ | $128^{2}$ | $192^{2}$ | $384^{2}$ |
| No-Prec | $201(0.56)$ | $1028(12)$ | $5170(316)$ | - | - |
| ILU $(A, 0)$ | $55(0.36)$ | $348(9)$ | $1484(131)$ | $2344(498)$ | - |
| ILU $(A, 1)$ | $26(0.14)$ | $126(4)$ | $577(62)$ | $894(207)$ | - |
| ILU $(M, 0)$ | $57(0.29)$ | $213(8)$ | $1289(122)$ | $2072(451)$ | - |
| ILU $(M, 1)$ | $28(0.28)$ | $116(4)$ | $443(48)$ | $763(191)$ | $2021(1875)$ |
| MG(V(1,1)) | $13(0.21)$ | $38(3)$ | $94(28)$ | $115(82)$ | $252(850)$ |

## Shifted Laplace Preconditioner Spectrum

- Eigenvalues of the preconditioned operator are bounded by 1
- Small eigenvalues move to zero, as $k$ increases.

Spectrum of $M^{-1} A$, where $\left(\beta_{1}, \beta_{2}\right)=(1,0.5)$

$$
k=30 \quad k=120
$$




## 3. Second Level Preconditioners

Number of GMRES iterations. Shifts in the preconditioner are $(1,0.5)$

| Grid | $k=10$ | $k=20$ | $k=30$ | $k=40$ | $k=50$ | $k=100$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $n=32$ | $5 / 10$ | $8 / 17$ | $14 / 28$ | $26 / 44$ | $42 / 70$ | $13 / 14$ |
| $n=64$ | $4 / 10$ | $6 / 17$ | $8 / 28$ | $12 / 36$ | $18 / 45$ | $173 / 163$ |
| $n=96$ | $3 / 10$ | $5 / 17$ | $7 / 27$ | $9 / 35$ | $12 / 43$ | $36 / 97$ |
| $n=128$ | $3 / 10$ | $4 / 17$ | $6 / 27$ | $7 / 35$ | $9 / 43$ | $36 / 85$ |
| $n=160$ | $3 / 10$ | $4 / 17$ | $5 / 27$ | $6 / 35$ | $8 / 43$ | $25 / 82$ |
| $n=320$ | $3 / 10$ | $4 / 17$ | $4 / 27$ | $5 / 35$ | $5 / 42$ | $10 / 80$ |

Erlangga and Nabben, 2008, seems to be independent of $k$.
with / without deflation.

## Deflation: or two-grid method

Deflation, a projection preconditioner

$$
P=I-A Q, \quad \text { with } \quad Q=Z E^{-1} Z^{T} \text { and } E=Z^{T} A Z
$$

where,
$Z \in R^{n \times r}$, with deflation vectors $Z=\left[z_{1}, \ldots, z_{r}\right], \operatorname{rank}(Z)=r \leq n$
Along with a traditional preconditioner $M$, deflated preconditioned system reads

$$
P M^{-1} A u=P M^{-1} g .
$$

Deflation vectors shifted the eigenvalues to zero.

## Spectrum as function of $k$



## Deflation for Helmholtz

With choice of multigrid inter-grid transfer operator (Prolongation) as deflation matrix, i.e. $Z=I_{h}^{2 h}$ and $Z^{T}=I_{2 h}^{h}$ then

$$
P_{h}=I_{h}-A_{h} Q_{h}, \quad \text { with } \quad Q_{h}=I_{h}^{2 h} A_{2 h}^{-1} I_{2 h}^{h} \quad \text { and } A_{2 h}=I_{2 h}^{h} A_{h} I_{h}^{2 h}
$$

where
$P_{h}$ can be interpreted as a coarse grid correction and
$Q_{h}$ as the coarse grid operator

## Deflation: ADEF1

Deflation can be implemented combined with SLP $M_{h}$,

$$
M_{h}^{-1} P_{h} A_{h} u_{h}=M_{h}^{-1} P_{h} g_{h}
$$

$A_{h} u_{h}=g_{h}$ is preconditioned by two-level preconditoner $M_{h}^{-1} P_{h}$.
For large problems, $A_{2 h}$ is too large to invert exactly. Inversion of $A_{2 h}$ is sensitive, since $P_{h}$ deflates the spectrum to zero.

To do is: Solve $A_{2 h}$ iteratively to required accuracy on certain levels, and shift the deflated spectrum to $\lambda_{h}^{\max }$ by adding a shift in two level preconditioner. This leads to the ADEF1 preconditioner

$$
P_{(h, A D E F 1)}=M_{h}^{-1} P_{h}+\lambda_{h}^{\max } Q_{h}
$$

## Deflation: MLKM

Multi Level Krylov Method a, take $\hat{A}_{h}=M_{h}^{-1} A_{h}$, and define $\hat{P}_{h}$ by using $\hat{A}_{h}$ (instead of $A_{h}$ ) will be

$$
\hat{P}_{h}=I_{h}-\hat{A}_{h} \hat{Q}_{h},
$$

where

$$
\hat{Q}_{h}=I_{h}^{2 h} \hat{A}_{2 h}^{-1} I_{2 h}^{h} \text { and } \hat{A}_{2 h}=I_{2 h}^{h} \hat{A}_{h} I_{h}^{2 h}=I_{2 h}^{h}\left(M_{h}^{-1} A_{h}\right) I_{h}^{2 h}
$$

Construction of coarse matrix $A_{2 h}$ at level $2 h$ costs inversion of preconditioner at level $h$. Approximate $A_{2 h}$

Ideal

$$
A_{2 h}=I_{2 h}^{h}\left(M_{h}^{-1} A_{h}\right) I_{h}^{2 h}
$$

## Practical

$$
\begin{aligned}
A_{2 h} & =I_{2 h}^{h}\left(M_{h}^{-1} A_{h}\right) I_{h}^{2 h} \\
A_{2 h} & \approx I_{2 h}^{h} I_{h}^{2 h} M_{2 h}^{-1} A_{2 h}
\end{aligned}
$$

## 4. Fourier Analysis

Dirichlet boundary conditions for analysis.
With above deflation,

$$
\mathbf{\operatorname { s p e c }}\left(P M^{-1} A\right)=f\left(\beta_{1}, \beta_{2}, k, h\right)
$$

is a complex valued function.
Setting $k h=0.625$,

- Spectrum of $P M^{-1} A$ with shifts $(1,0.5)$ is clustered around 1 with a few outliers.
- Spectrum remains almost the same, when the imaginary shift for the preconditioner is varied from 0.5 to 1 .


## Fourier Analysis

ADEF1: Analysis shows spectrum clustered around 1 with few outliers.

$$
k=30
$$

$k=120$



## Fourier Analysis

Spectrum of Helmholtz preconditioned by MLKM ${ }^{\oplus}$,
$k=160$ and $20 \mathrm{gp} / \mathrm{wl}$

Ideal


Practical


## 5. Numerical results



Number of GMRES iterations for the 1D Helmholtz equation

$$
10 \leq k \leq 800
$$

## Numerical results



Number of GMRES iterations for the 1D Helmholtz equation

$$
1000 \leq k \leq 20000
$$

## Numerical results

Number of GMRES outer-iterations in multilevel algorithm.
$\left(\beta_{1}, \beta_{2}\right)=(1,0.5) k h=.3125$ or $20 \mathrm{gp} / \mathrm{wl}$ and SLP approximated by multigrid Vcyle $\mathrm{V}(1,1)$

| Grid | $k=10$ | $k=20$ | $k=40$ | $k=80$ | $k=160$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| ADEF1-V(4,2,1) | 9 | 11 | 16 | 27 | $100+$ |
| ADEF1-V(6,2,1) | 9 | 10 | 14 | 21 | 47 |
| ADEF1-V(8,2,1) | 9 | 10 | 13 | 20 | 38 |
| ADEF1-V(8,3,2) | 9 | 10 | 13 | 19 | 37 |

ADEF1-V(8,2,1), a multilevel solver where 8 and 2 Krylv iterations performed on $2^{\text {nd }}$ and $3^{\text {rd }}$ levels and 1 iteration on further levels and Vcycle approximates SLP.

## Results

Petsc solve-time in Seconds; a Two-level solver.

| Solver | 20 | 40 | 80 | 120 | 160 | 200 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| SLP | $0.01(23)$ | $0.24(54)$ | $2.62(113)$ | $11.60(168)$ | $33.59(222)$ | $83.67(274)$ |
| ADEF1/SLP | $0.03(10)$ | $0.14(14)$ | $0.82(23)$ | $2.92(37)$ | $8.98(61)$ | $23.13(87)$ |

SLP : GCR preconditioned with SLP $M(1,1)$.
Def/SLP: Deflated and preconditioned GCR.
Grid resolution is such that there are 10 grid points per wavelength.

## Numerical results



Comparison of number of iterations by ADEF1 and MLKM.

## Numerical results

3D Helmholtz on unit cube with Sommerfeld b.c. on all faces.

|  | Wavenumber $k$ |  |  |  |  |  |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Solver Type | 5 | 10 | 15 | 20 | 30 | 40 | 60 | 80 |
| SL Prec. | 11 | 15 | 21 | 29 | 47 | 74 | 118 | 185 |
| ADEF1-F(8,2,1) | 9 | 10 | 11 | 11 | 13 | 16 | 22 | 28 |
| ADEF1-V(8,2,1) | 11 | 15 | 21 | 28 | 44 | 67 | 101 | 153 |

SL Prec. : Only shifted Laplace preconditioner
ADEF1-F : Multilevel solver , Fcycle for slp.

ADEF1-V : Multilevel solver, Vcycle for slp.

## Numerical results

Multilevel solver is coded in Petsc.
3D Helmholtz on unit cube with constant wavenumber.



ADEF1 solve time and Setup time.

## Conclusions

- Without deflation, when imaginary shift is increased in SLP, spectrum remains bounded above 1, but lower part moves to zero.
- With deflation the convergence is nearly independent of the imaginary shift.
- With deflation the convergence is initially weakly depending on $k$. For large $k$ is scales again linearly.
- With deflation the CPU time is less than without deflation.
- The convergence of ADEF1 and the practical variant of MLKM are similar.


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